

braneworld gravity

Joseph Lykken
Fermilab

collaborators: Ruoyu Bao, Marcela Carena, Minjoon Park, Jose Santiago

papers: [hep-ph/0506305](#) + ...

outline

- gravity in various 5d braneworld orbifolds
- problems with DGP
- can we make gravity decouple?
- the interval approach to braneworld gravity
- straight gauges
- some general results
- ghosts, strong coupling, and vDVZ

what is missing from this talk

- codimension > 1 braneworlds
- anything beyond pure gravity
- dynamical branes, strings
- SUSY
- proper attribution to all the brilliant+charming people who have contributed to this subject

a brane in 5d



$$2M^3 \int d^5x \sqrt{-G} R + \int d^4x^0 \sqrt{-g^0} (\mathcal{L}_m - 2M^3 U_0)$$

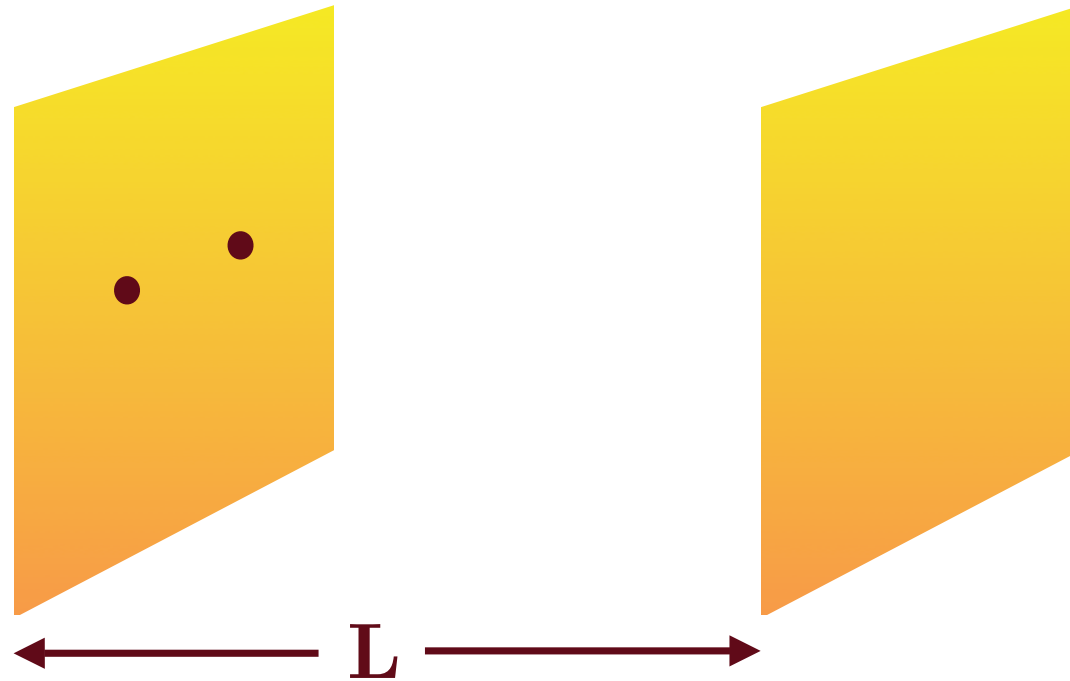
brane tension $U_0 = 0 \Rightarrow$ flat 5d background

gravitational potential between
two brane sources:

$$\propto \frac{1}{M^3 r^2}$$

gravity on an orbifold

same as previous case, but compactify the 5th dimension on a Z_2 orbifold of a circle of circumference $2L$



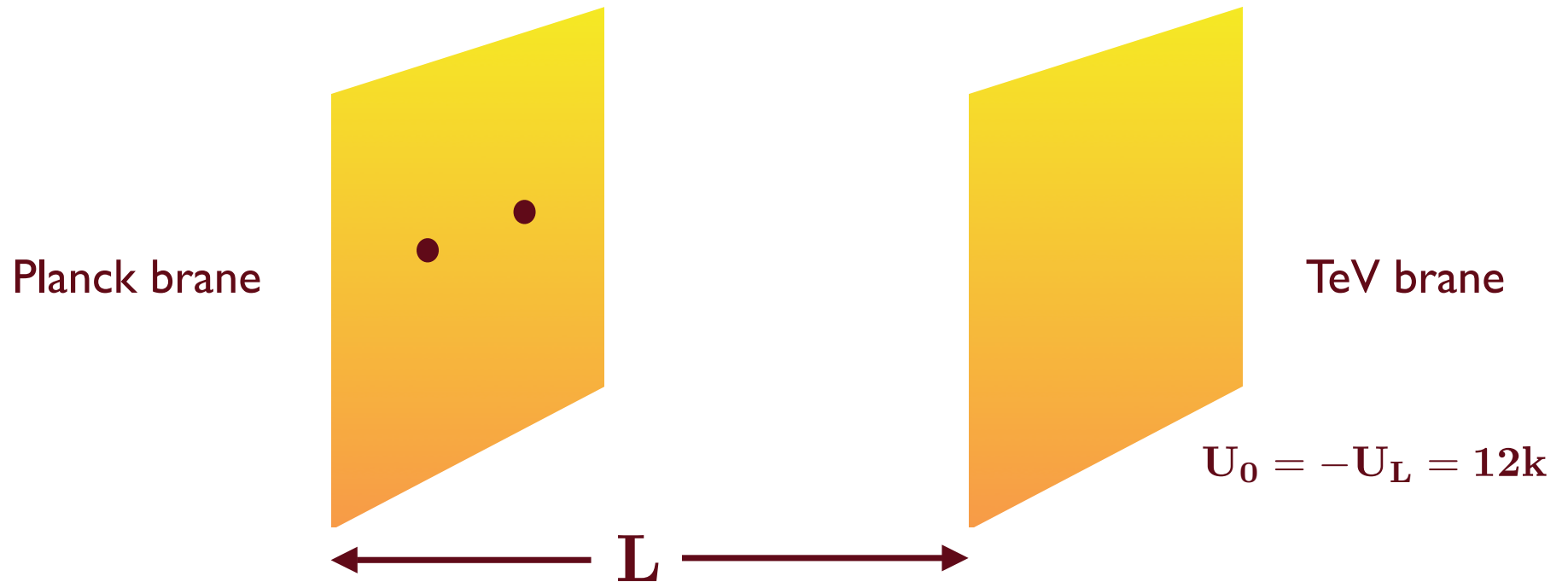
orbifolding means that graviton KK modes obey Neumann b.c. at both branes

gravitational potential between two brane sources:

$$\propto \frac{1}{M_4^2 r} + \sum_n \frac{1}{M_4^2 r} e^{-\frac{nr}{L}} \quad \text{where } M_4^2 \equiv M^3 L$$

RSI model

same as previous case, but add a bulk AdS cosm. constant $\Lambda = -24M^3k$ where $k \simeq M$ and tune the brane tensions to make 4d slices flat

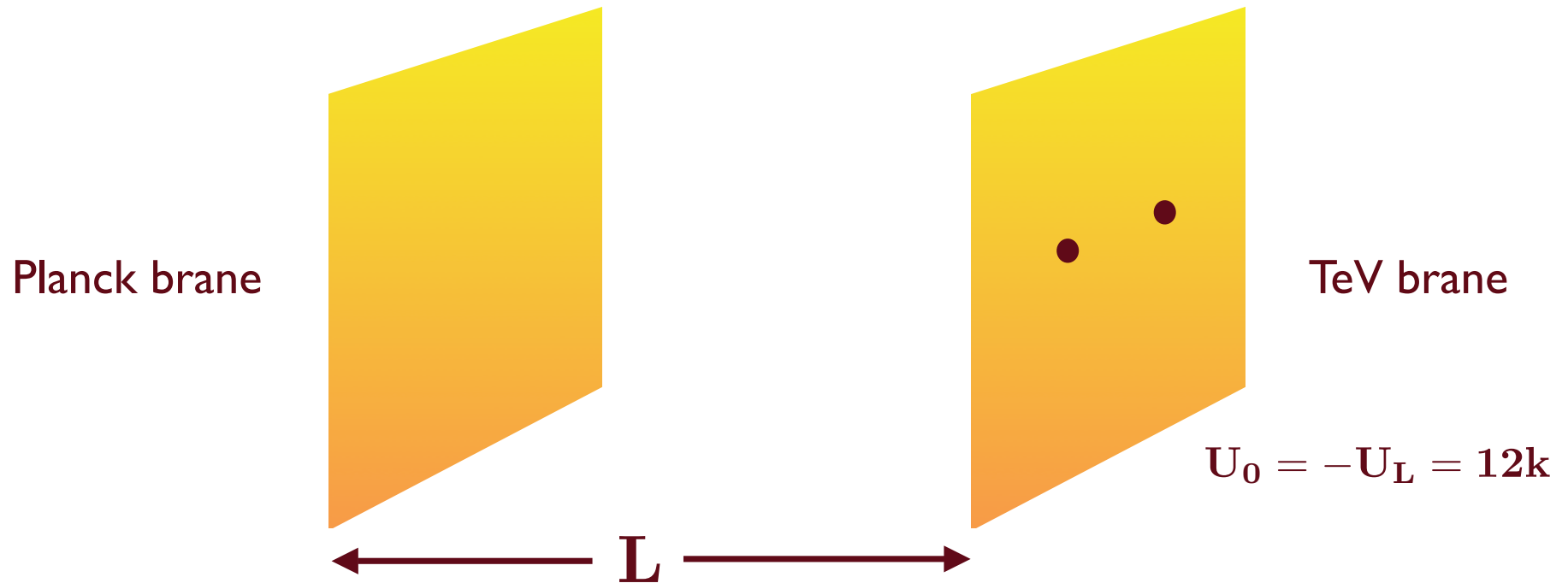


there is a normalizable graviton zero mode + a tower of massive KK modes

gravitational potential between two Planck brane sources: $\propto \frac{1}{M_4^2 r} + \frac{1}{M_4^2 k^2 r^3}$ where $M_4^2 \simeq M^3/k$

RSI model

due to warping, at the TeV brane $M, k \rightarrow T, T = \text{TeV}$



but the graviton zero mode exchange also has a wave function suppression $\sim T^2/k^2$

gravitational potential between two TeV brane sources:

$$\propto \frac{1}{M_4^2 r} + \frac{1}{T^4 r^3} \quad \text{LHC!}$$

RSII and LR models

take $L \rightarrow \infty$ limit, look at gravity on the Planck brane (RSII) or on a probe brane (LR)



due to warping, graviton zero mode is still normalizable.

there is a continuum of massive KK graviton modes

gravitational potential on the Planck brane is \sim same as in RS I

gravitational potential between two probe brane sources:

$$\propto \frac{1}{M_4^2 r} + \frac{1}{T^8 r^7} \quad \text{LHC!}$$

detuned warped models

Planck brane



TeV brane



← L →

gives $AdS_5 AdS_4$

for any $-12 \leq U_i \leq 12$

$$g_{\mu\nu} = \frac{a(y)^2}{\left(1 - \frac{H^2 x^2}{4}\right)^2} \eta_{\mu\nu}$$

warp factor: $a(y) = \frac{\cosh k(y - y_0)}{\cosh ky_0}$ where: $\tanh ky_0 = \frac{U_0}{12k}$

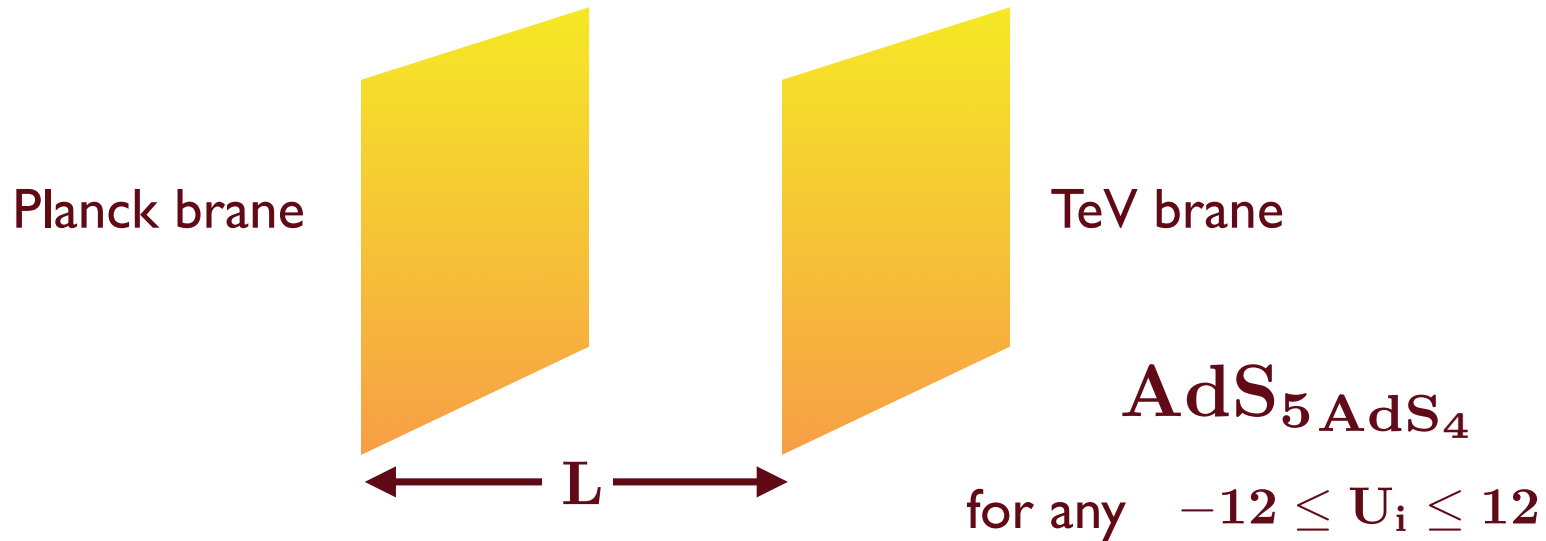
$\tanh k(L - y_0) = \frac{U_L}{12k}$

$H = \frac{k}{\cosh ky_0}$

- * L is now fixed
- * warp factor blows up for large y, instead of damping out as in RSII or LR
- * in the “bigravity” region, $U_0 = U_L < \simeq 12k$, there is **both** a zero mode graviton and an ultralight KK graviton, with mass $\sim H^2/k$

Kogan, Mouslopoulos, Papazoglou, Ross, Santiago, Kaloper, etc

detuned warped models



warp factor: $a(y) = \frac{\cosh k(y - y_0)}{\cosh ky_0}$ where: $\tanh ky_0 = \frac{U_0}{12k}$

$\tanh k(L - y_0) = \frac{U_L}{12k}$

$$H = \frac{k}{\cosh ky_0}$$

- * $L \rightarrow \infty$ removes the TeV brane, but also removes the graviton zero mode, which is no longer normalizable.
- * this leaves a model with only the ultralight KK mode to mock up 4d gravity on the Planck brane (Karch-Randall)

DGP model

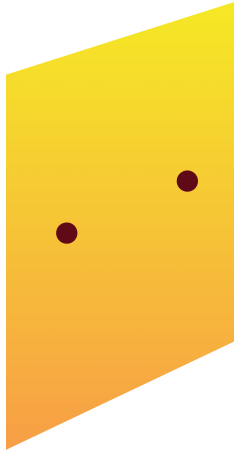


$$2M^3 \int d^5x \sqrt{-G} R + 2M_4^2 \int d^4x^0 \sqrt{-g^0} \tilde{R}$$

- a single tensionless brane ($U=0$)
- flat 5d Minkowski background
- add brane kinetic term for gravity
- the two Planck constants compete against each other

Dvali, Gabadadze, Porrati

DGP model



$$2M^3 \int d^5x \sqrt{-G} R + 2M_4^2 \int d^4x^0 \sqrt{-g^0} \tilde{R}$$

$$\lambda \equiv M_4^2 / M^3$$

graviton propagator: $\frac{P_{\mu\nu\rho\lambda}}{M^3 p + M_4^2 p^2}$

- the two Planck constants compete against each other
- for brane observer, gravity looks 5d on distance scales $> \lambda$; could self-accelerate the expansion of our (brane) universe!
- at shorter scales, it looks like 4d gravity + an extra scalar interaction

$$P_{\mu\nu\rho\lambda} = \frac{1}{2}(\eta_{\mu\rho}\eta_{\nu\lambda} + \eta_{\mu\lambda}\eta_{\nu\rho}) - \frac{1}{3}\eta_{\mu\nu}\eta_{\rho\lambda}$$

problems with the DGP model



graviton propagator: $\frac{\mathbf{P}_{\mu\nu\rho\lambda}}{M^3\mathbf{p} + M_4^2\mathbf{p}^2}$

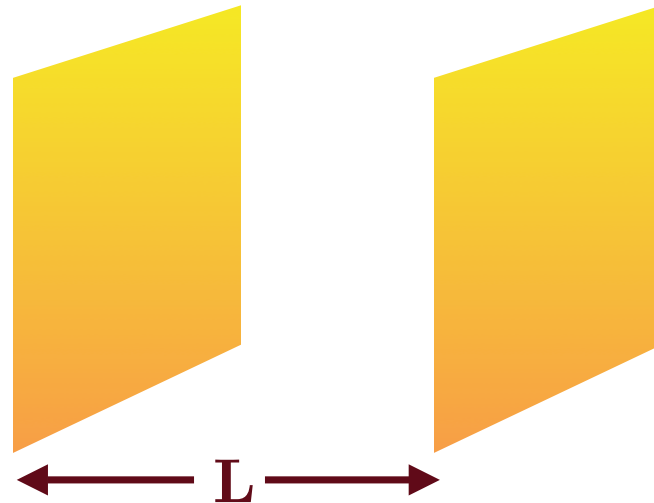
$$\mathbf{P}_{\mu\nu\rho\lambda}^{\text{DGP}} = \frac{1}{2}(\eta_{\mu\rho}\eta_{\nu\lambda} + \eta_{\mu\lambda}\eta_{\nu\rho}) - \frac{1}{3}\eta_{\mu\nu}\eta_{\rho\lambda}$$

$$\mathbf{P}_{\mu\nu\rho\lambda}^{4\text{d}} = \frac{1}{2}(\eta_{\mu\rho}\eta_{\nu\lambda} + \eta_{\mu\lambda}\eta_{\nu\rho}) - \frac{1}{2}\eta_{\mu\nu}\eta_{\rho\lambda}$$

- the “4d” graviton has the tensor structure of a 5d graviton; related to the vDVZ discontinuity: in flat space, the massless limit of a massive graviton is *not* a massless graviton
- there is an extra scalar mode whose kinetic term in the “brane effective action” vanishes as $\lambda = M_4^2/M^3 \rightarrow \infty$; this implies that strong coupling sets in at a (low) energy scale M^2/M_4 , linearized gravity breaks down, and the low energy effective field theory is inconsistent

Luty, Porrati, Rattazzi; Rubakov, etc

can we make gravity decouple?

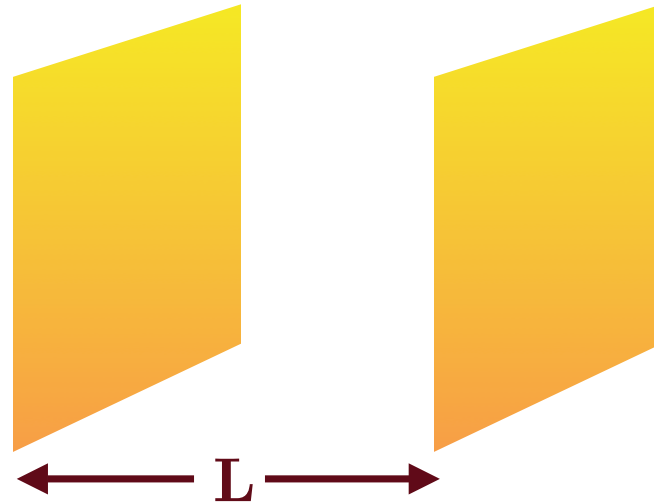


- first do it for gauge theory:
- go back to the Z_2 orbifold, warped or flat
- put an abelian gauge theory in the bulk
- add brane kinetic terms on both branes

$$\frac{1}{4g_5^2} \int d^5x F_{MN} F^{MN} + \sum_i \frac{1}{4g_4^{(i)2}} \int d^4x F_{\mu\nu}^{(i)} F^{\mu\nu (i)}$$

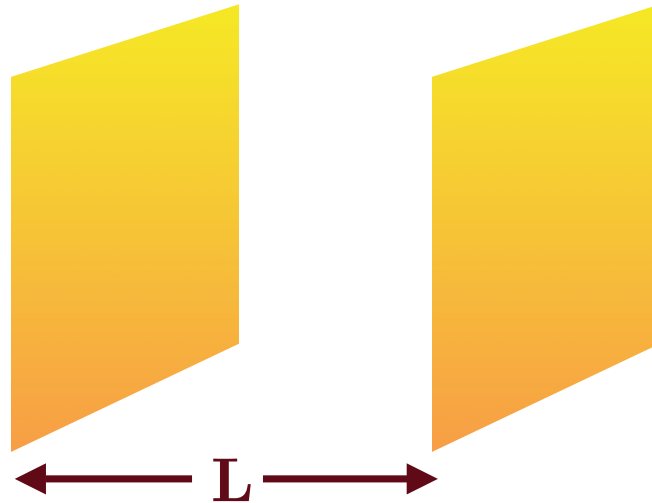
Carena, Ponton, Tait, Wagner
Davoudiasl, Hewett, Rizzo

decoupling gauge theories



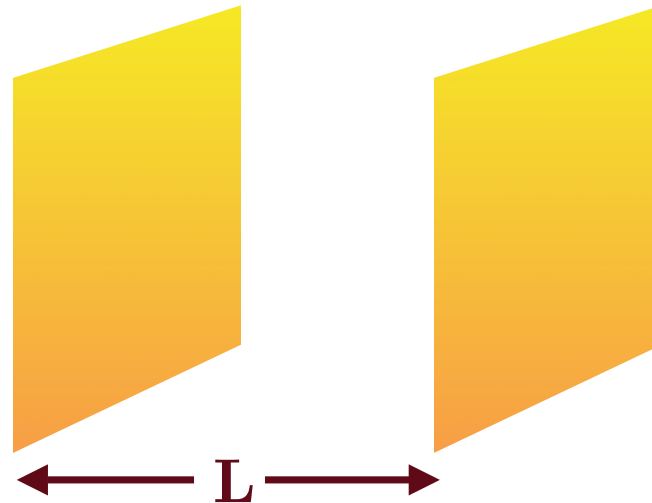
- in addition to the zero mode photon, there is an ultralight massive KK photon
- charged sources on the Planck brane couple to only one linear combination of these two modes
- charged sources on the TeV brane couple only to the orthogonal linear combination
- in a certain energy range there are two decoupled 4d gauge theories

decoupling gauge theories



- this is the beginning of deconstructing or latticizing 5d gauge theory
- there is no problem with having two photons in the effective 4d theory, since we can have two different gauge symmetries
- and there is no vDVZ discontinuity for photons

decoupling gravity?



- however we *do* expect problems in deconstructing or latticizing 5d gravity
- there is a big problem with having two *gravitons* in the effective 4d theory, since there is only one 4d general covariance
- and there is a vDVZ discontinuity for gravitons (in flat space)
- but if you *could* decouple a 4d graviton from (ordinary!) matter sources on the “wrong” brane, it would be a big deal

Arkani-Hamed, Georgi, Schwartz, Randall, etc

how to make progress on these issues?

- many worthy researchers have proposed many brilliant and novel approaches to this complex of problems/opportunities
- let's ignore all of them in favor of the most simple-minded approach imaginable:

a straight ahead approach

- write the action for the most general class of 5d warped brane gravity models that you have the patience to analyze
- gauge-fix them
- integrate out the extra dimension to get the gauge-fixed 4d effective gravity theory
- in the big parameter space, try to generalize potentially interesting “limits”
- in the big parameter space, explore issues of strong coupling, vDVZ, decoupling, and their relation e.g. to kinetic ghosts (an older but better-understood problem)

why is it hard?

- write the action for the most general class of 5d warped brane gravity models: need the “interval picture” for braneworld orbifold gravity
- gauge-fix them: need to invent “straight gauges”
- integrate out the extra dimension to get the gauge-fixed 4d effective gravity theory: nasty calculations, especially if you want to go beyond quadratic order to examine strong coupling issues

gravity on orbifolds?

- general relativity is defined on manifolds, not orbifolds; orbifolds have singular fixed points
- this is not a problem in string theory, because the strings don't care about the singularities, and anyway they are just limits of smooth manifolds in the bigger moduli space
- does this mean we have to “blow up” the fixed points of our orbifold braneworlds, in order to have a well-defined gravity action?

the interval approach to braneworld gravity

- replace the 5d orbifold by two 5d manifolds with boundaries; the branes are wedged “in between” adjacent boundaries
- a general 5d warped gravity model is then described by an action like

$$\mathbf{S} = \left(\int_{\mathcal{M}_1} d^5\mathbf{x} + \int_{\mathcal{M}_2} d^5\mathbf{x} \right) \sqrt{-\mathbf{G}} \left(2\mathbf{M}^3 \mathbf{R} - \Lambda \right) \\ + 2\mathbf{M}^3 \sum_{\mathbf{i}} \int_{\Phi_{\mathbf{i}}=0} d^4\mathbf{x}^{(\mathbf{i})} \sqrt{-\mathbf{g}^{(\mathbf{i})}} \left(\lambda_{\mathbf{i}} \tilde{\mathcal{R}}^{(\mathbf{i})} - \mathbf{U}_{\mathbf{i}} \right) + 4\mathbf{M}^3 \oint_{\partial\mathcal{M}_1 + \partial\mathcal{M}_2} \mathbf{K}$$

brane-boundary equations

- this representation of the gravity orbifold has a well-defined action principle
- but the bulk Einstein equations have supplemented by additional “brane-boundary” equations:

$$\text{bulk: } \mathbf{R}_{MN} - \frac{\mathbf{G}_{MN}}{2} \left(\mathbf{R} - \frac{\Lambda}{2M^3} \right) = 0$$

brane-boundary:

$$\left[\left(\lambda_i \tilde{\mathcal{R}}_{\alpha\beta}^{(i)} - \frac{\mathbf{g}_{\alpha\beta}^{(i)}}{2} (\lambda_i \tilde{\mathcal{R}}^{(i)} - \mathbf{U}_i) \right) \mathbf{e}^{(i)\bar{\alpha}}_{\mathbf{M}} \mathbf{e}^{(i)\bar{\beta}}_{\mathbf{N}} + 2\mathbf{P}^{(i)}_{\mathbf{M}} \nabla_{\mathbf{P}} \mathbf{N}^{(i)}_{\mathbf{N}} - 2\mathbf{P}^{(i)}_{\mathbf{MN}} \nabla_{\mathbf{P}} \mathbf{N}^{(i)\mathbf{P}} \right]_{\Phi_i=0} = 0$$

where $\mathbf{N}^{(i)}_{\mathbf{M}}$ is the outward normal to the i th brane/boundary and $\mathbf{P}^{(i)}_{\mathbf{MN}}$ is the projection operator onto the i th brane/boundary

a problem with gauge-fixing

- to get the effective 4d action for the physical modes, we need to do a complete gauge fixing of the 5d general coordinate invariance
- the usual gauge choices (harmonic, de Donder, axial) that simplify the equations of motion, have a big problem:
- in these gauges at least one of the branes looks “bent”

straight gauges

- possible solutions are: (1) don't try to do a complete gauge-fixing, (2) analyze the 4d theory with bent branes, and (3) analyze the 4d theory in two different coordinate patches
- instead we have defined a family of “straight gauges”
- straight gauge = a choice of 5d bulk coordinate system such that:
 - both of the branes are described by straight slices $y = y_i$
 - $\left[\mathbf{G}_{\mu 4}^{(i)} \right]_{y=y_i} = 0$

straight gauges

- good thing about straight gauges: the branes are straight, and the brane-boundary equations simplify a lot:

$$\left[\lambda_i \tilde{\mathcal{R}}_{\mu\nu} - \frac{1}{2} \mathbf{g}_{\mu\nu} (\lambda_i \tilde{\mathcal{R}} - \mathbf{U}_i) + \theta_i \sqrt{\mathbf{G}^{44}} (\mathbf{g}'_{\mu\nu} - \mathbf{g}_{\mu\nu} \mathbf{g}'_{\rho\sigma} \mathbf{g}^{\rho\sigma}) \right]_{\mathbf{y}=\mathbf{y}_i} = \mathbf{0}$$

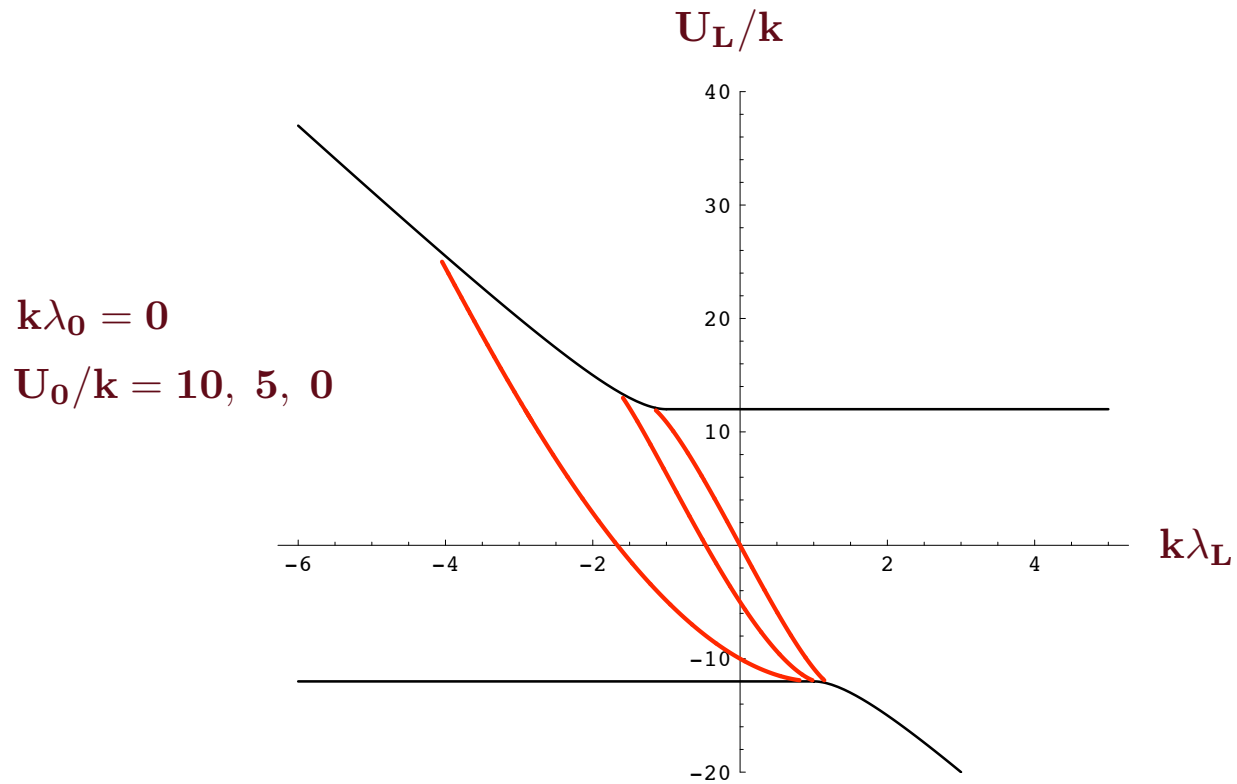
- bad thing about straight gauges: the full analysis of the equations of motion is tedious

some general results

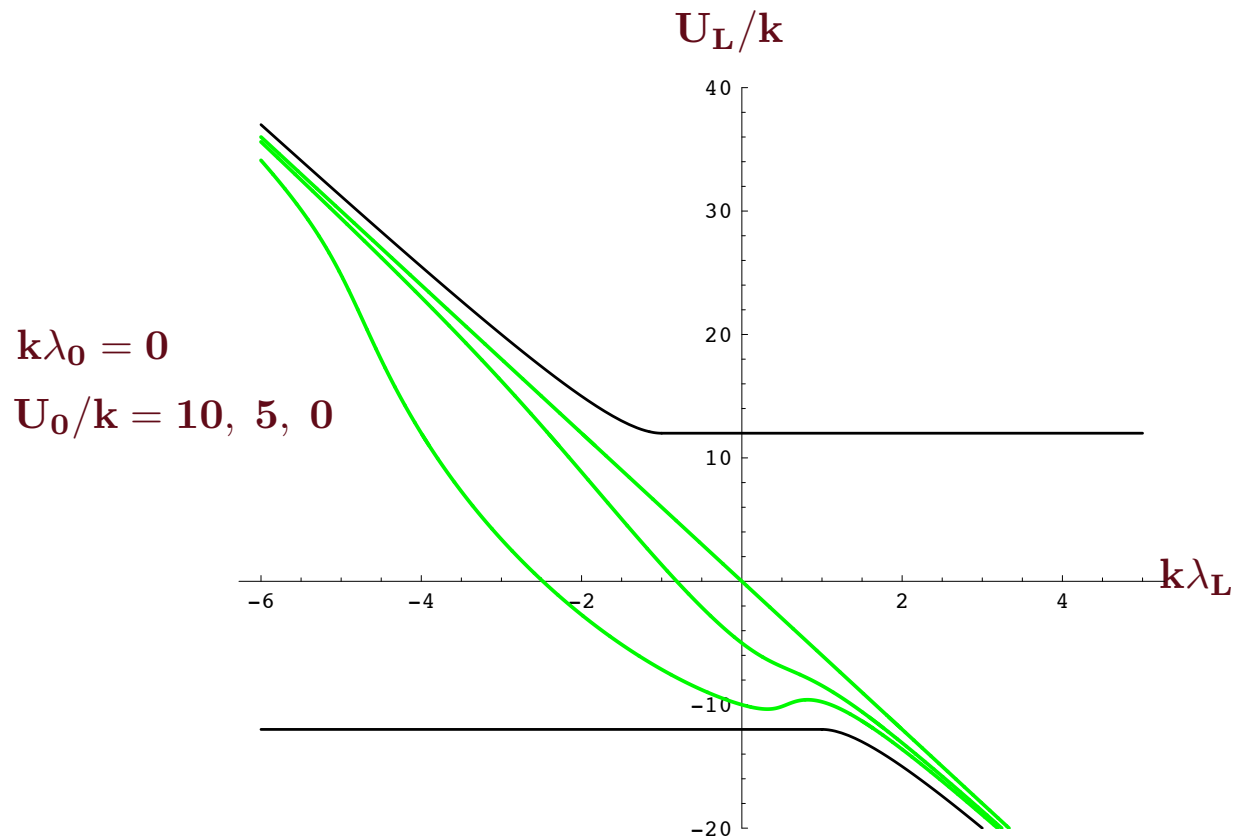
- looked at most general interval picture action for 5d braneworld gravity, with two branes, up to second order in derivatives
- expanded around a general AdS_5AdS_4 background geometry
- gauge-fixed to a straight gauge, and computed the quadratic 4d effective action for the physical modes
- the only physical modes are the graviton and a radion

ghosts, strong coupling, vDVZ

- there are 4 input parameters which define the AdS_5AdS_4 background geometry: two brane tensions U_0, U_L , and the coefficients of the brane kinetic terms λ_0, λ_L
- we compute the coefficients in the 4d effective action of the kinetic terms for the radion and for the zero mode of the graviton
- in each case, these coefficients can be positive (that's good), negative (that's a kinetic ghost), or zero (that's strong coupling).
- we can also check for vDVZ problems in various limits



- inside the black lines, there is an $AdS_5 AdS_4$ background solution
- to the left of the red lines, the zero mode graviton is a kinetic ghost
- at the red lines, there is strong coupling



- inside the black lines, there is an $AdS_5 AdS_4$ background solution
- to the left of the green lines, the radion is a kinetic ghost
- at the green lines, there is strong coupling

is this general enough to include DGP-like limits?

- yes, but this is work in progress
- note that the DGP-like limits are living in the same parameter space as lots of other models, including RS, KR, and bigravity.

outlook

- we still have a lot to learn about braneworld models that mock up 4d gravity over some range of length/energy scales
- this is an important topic for both LHC and cosmology

