SUSY Precision Physics
at Colliders

W. Hollik, MPI München

SUSY 2005
Durham, July 18 - 23, 2005

- Introduction
- Precision physics with standard particles
- Precision physics with Higgs bosons
- Precision physics with SUSY particles

Precision analysis required for

- Indirect tests of the MSSM
  → virtual SUSY effects in precision observables
- Precision studies for SUSY particles
  → determination of masses & couplings
  → reconstruction of model parameters
- Direct versus indirect tests
  → precision observables for precisely measured SUSY parameters
  → consistency check

Processes with external

(i) standard particles

(ii) Higgs bosons, especially light Higgs h^0

(iii) SUSY particles
– the chargino and neutralino sector
– the sfermion sector

(expected) experimental precision

error for	LEP/Tev	Tev/LHC	LC	GigaZ
M_W [MeV]	33	15	15	7
$\sin^2\theta_{\text{eff}}$	0.00017	0.00021		0.000013
m_{top} [GeV]	4.3	2	0.2	0.13
M_{Higgs} [GeV]	–	0.1	0.05	0.05

together with

$\delta M_Z = 2.1$ MeV (LEP)

$\delta G_{\text{F}}/G_{\text{F}} = 1 \cdot 10^{-5}$ (μ lifetime)

Detailed analysis for SPS1a benchmark scenario: potential of LHC (300 fb^{-1}) alone and LHC $+$ LC

	LHC	LHC+LC
$\Delta m_{\tilde{\chi}_1^0}$	4.8	0.05 (input)
$\Delta m_{\tilde{l}_R}$	4.8	0.05 (input)
$\Delta m_{\tilde{\chi}_2^0}$	4.7	0.08
$\Delta m_{\tilde{q}_L}$	8.7	4.9
$\Delta m_{\tilde{q}_R}$	11.8	10.9
$\Delta m_{\tilde{g}}$	8.0	6.4
$\Delta m_{\tilde{b}_1}$	7.5	5.7
$\Delta m_{\tilde{b}_2}$	7.9	6.2
$\Delta m_{\tilde{l}_L}$	5.0	0.2 (input)
$\Delta m_{\tilde{\chi}_4^0}$	5.1	2.23

LHC+LC accuracy limited by LHC jet energy scale resolution

SPS 1a benchmark scenario:

favorable scenario for both LHC and LC

$\Rightarrow$ LC input improves accuracy significantly

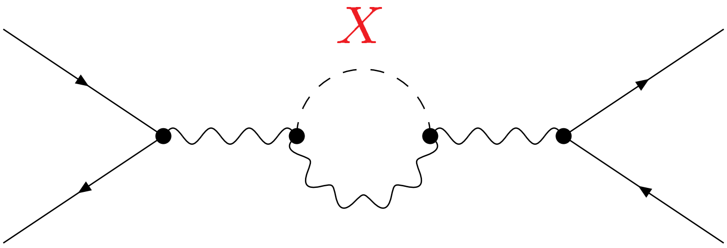
Comparison of electro-weak precision observables with theory:

EW Precision data:		Theory:
$\Delta\rho, \Delta r, M_W, \sin^2\theta_{\text{eff}}$	↔	SM, MSSM , . . .

$\Downarrow$

Test of theory at quantum level:

Sensitivity to loop corrections



$\Downarrow$

sensitivity to internal particles (X)

Precision observables: M_W, $\sin^2\theta_{\rm eff}$, m_h, $(g-2)_\mu$, b physics, . . .

1.) Theoretical prediction for M_W in terms of $M_Z, \alpha, G_\mu, \Delta r$:

$$M_W^2\left(1-\frac{M_W^2}{M_Z^2}\right)=\frac{\pi\,\alpha}{\sqrt{2}\,G_\mu}\left(\frac{1}{1-\Delta r}\right)$$

$\Updownarrow$

loop corrections

2.) Effective mixing angle:

$$\sin^2\theta_{\rm eff}=\frac{1}{4\,|Q_f|}\left(1-\frac{\mathrm{Re}\,g_V^f}{\mathrm{Re}\,g_A^f}\right)$$

Higher order contributions:

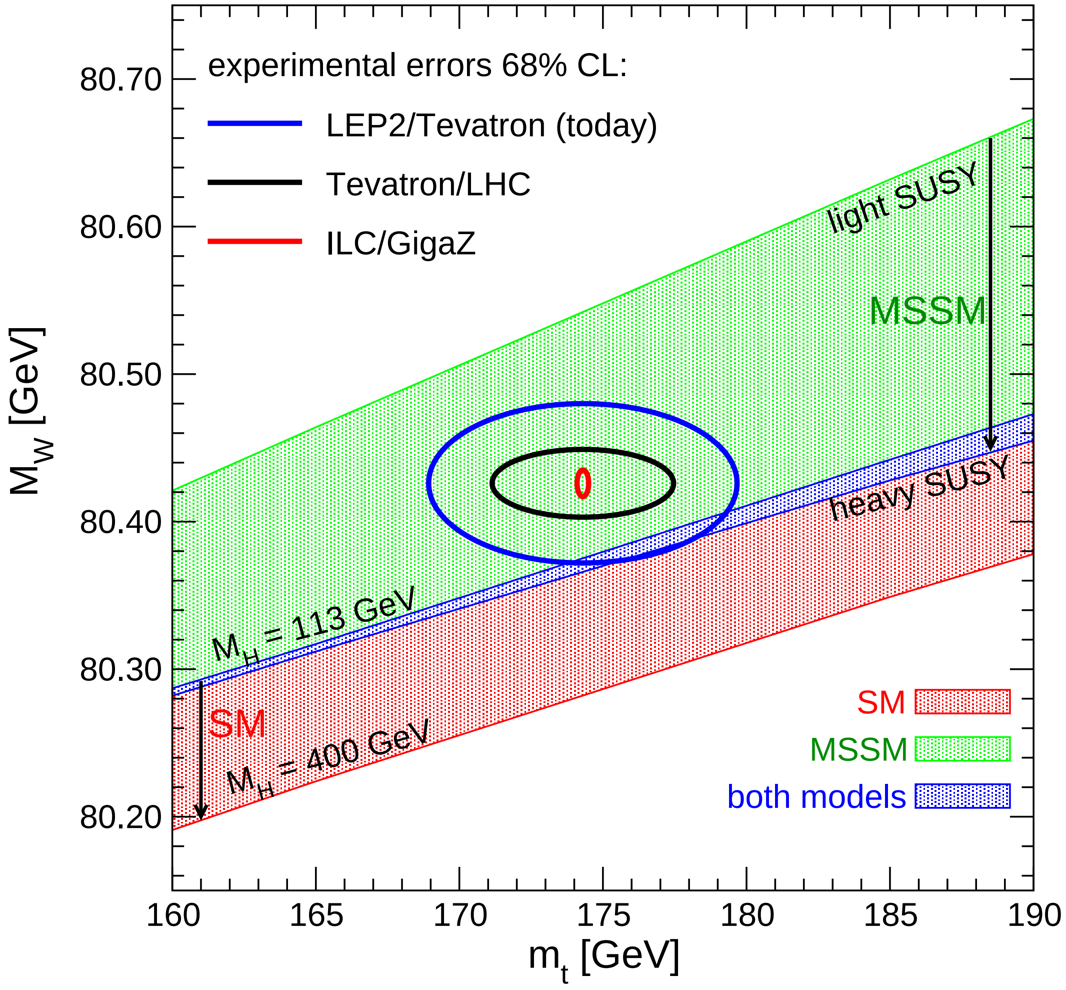
$$g_V^f \to g_V^f + \Delta g_V^f, \quad g_A^f \to g_A^f + \Delta g_A^f$$

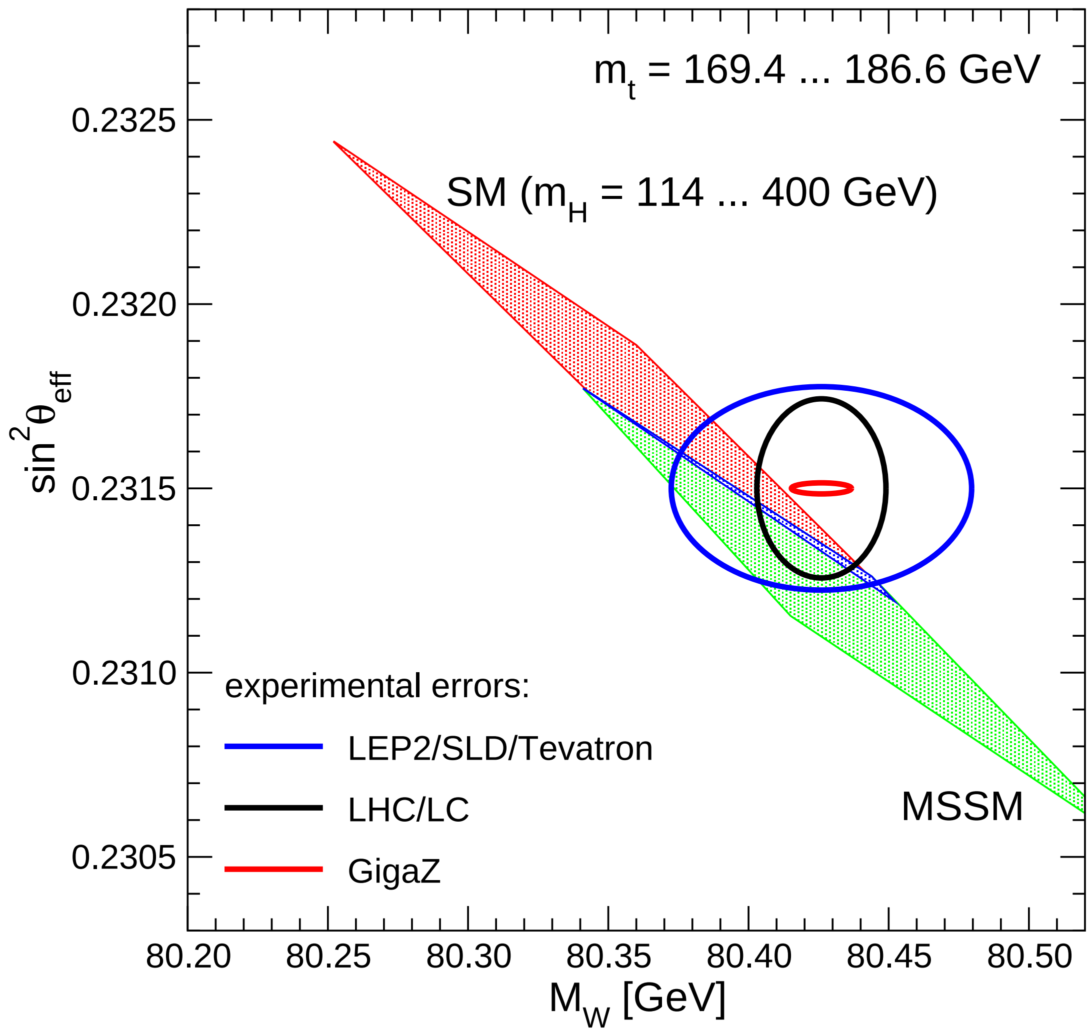
recent review:

Heinemeyer, WH, Weiglein, hep-ph/0412214

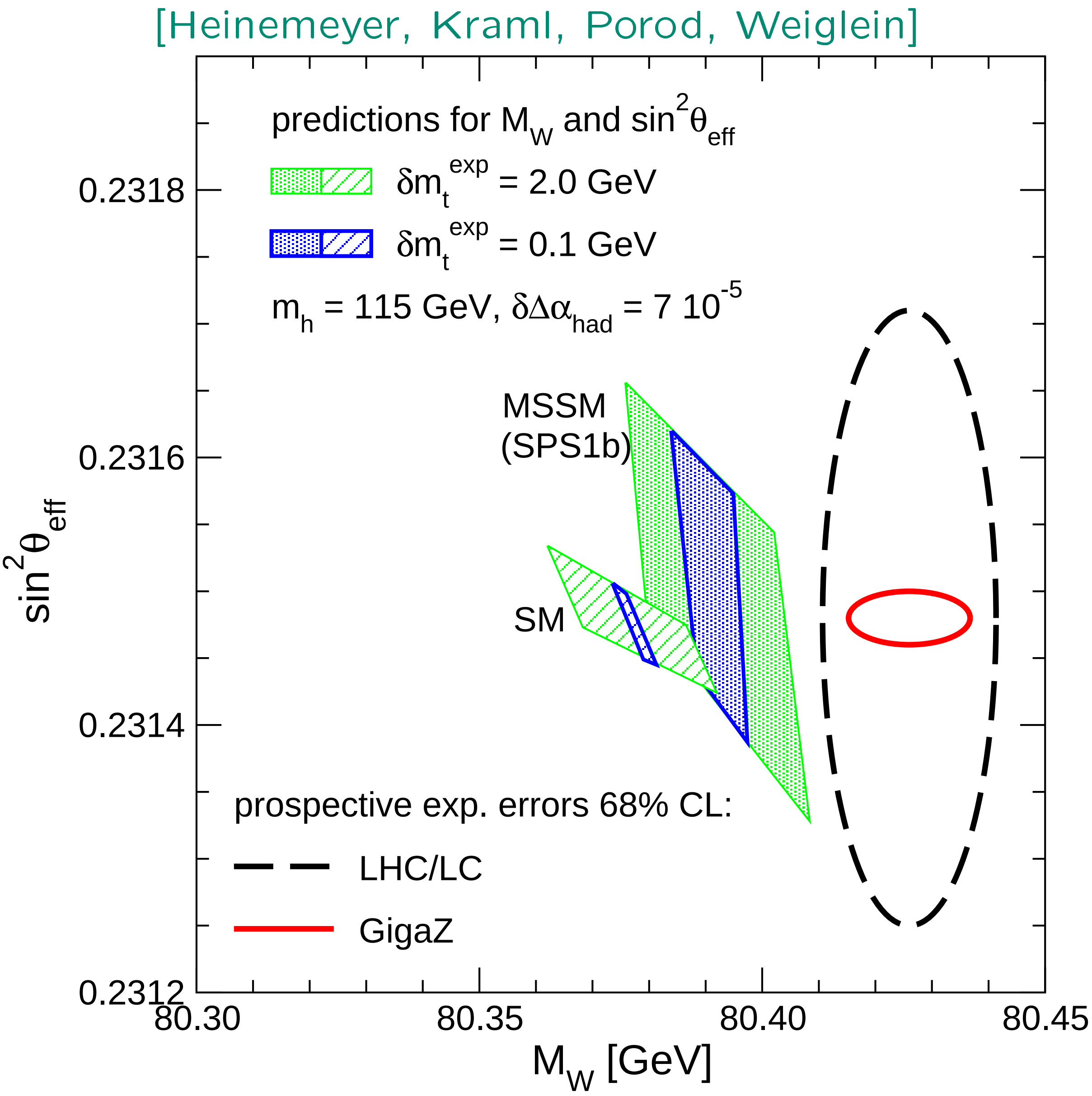
[Chankowski, Dabelstein, WH, Mösle, Pokorski, Rosiek]

[update: Heinemeyer, Weiglein]





m_t = 169.4 ... 186.6 GeV
SM (m_H = 114 ... 400 GeV)
MSSM
experimental errors:
LEP2/SLD/Tevatron
LHC/LC
GigaZ
0.2325
0.2320
0.2315
0.2310
0.2305
$\sin^2\theta_{\mathrm{eff}}$
80.20
80.25
80.30
80.35
80.40
80.45
80.50
M_W [GeV]

[Heinemeyer, Kraml, Porod, Weiglein]
predictions for M_W and $\sin^2\theta_{eff}$
δm_t^{exp} = 2.0 GeV
δm_t^{exp} = 0.1 GeV
m_h = 115 GeV, $\delta\Delta\alpha_{had}$ = 7 10^{-5}
MSSM
(SPS1b)
SM
prospective exp. errors 68% CL:
LHC/LC
GigaZ
0.2318
0.2316
0.2314
0.2312
80.30
80.35
80.40
80.45
$\sin^2\theta_{eff}$
M_W [GeV]

Models of SUSY breaking

generic MSSM: 105 parameters (masses, mixing angles, phases)

reduced to few parameters in specific models

mSUGRA: $m_0, m_{1/2}, A_0, \tan\beta, \mathrm{sign}(\mu)$

GMSB: $M_{\mathrm{mess}}, N_{\mathrm{mess}}, \tan\beta, \mathrm{sign}(\mu)$

AMSB: $m_{\mathrm{aux}}, m_0, \tan\beta, \mathrm{sign}(\mu)$

→ mass parameters at the electroweak scale
$(M_1, M_2, M_3, \mu, M_{\tilde{f}_{L,R}}, \ldots)$

Benchmark scenarios

"Snowmass points and slopes" (SPS),
hep-ph/0202233

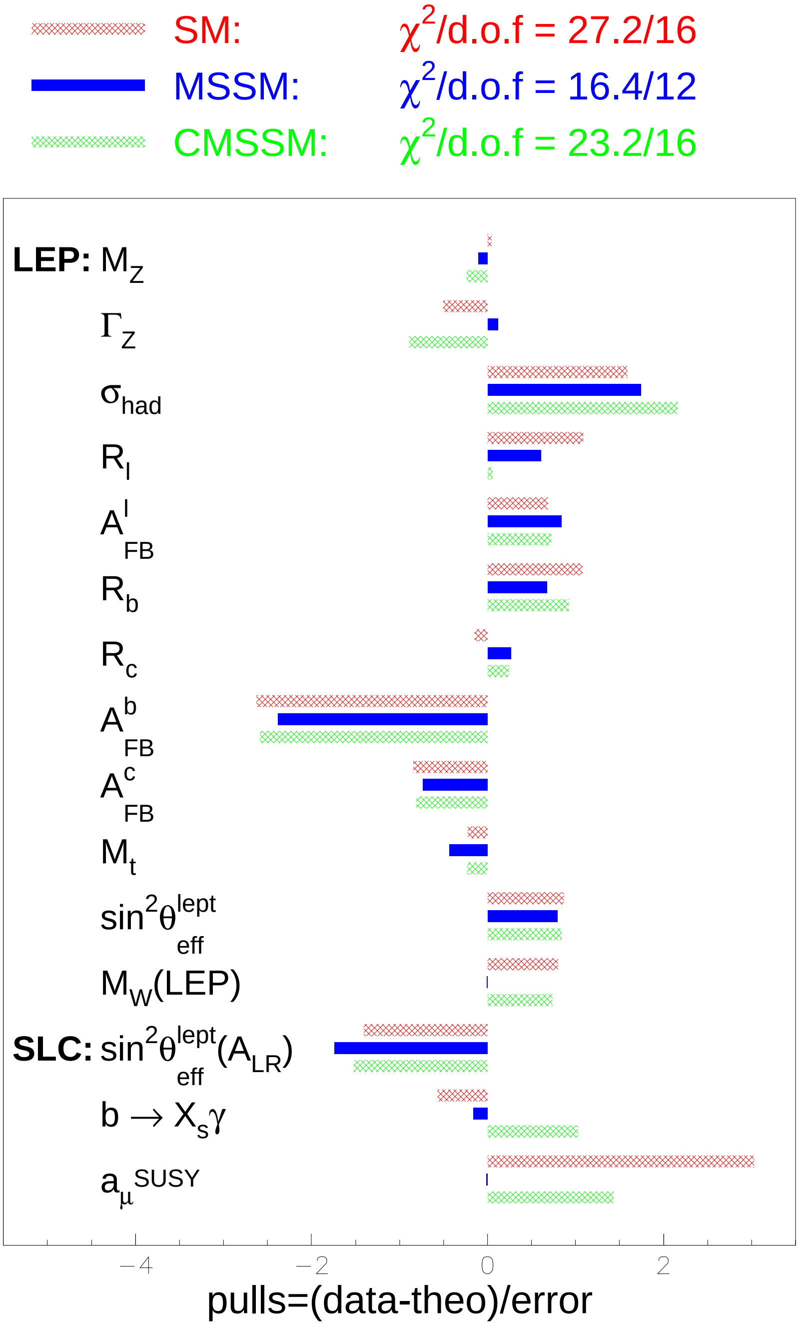
examples (mSUGRA):

- SPS1a: $m_0 = 100$ GeV, $m_{1/2} = 250$ GeV, $A_0 = -100$,
  $\tan\beta = 10$, $\mu > 0$.
- SPS1b: $m_0 = 200$ GeV, $m_{1/2} = 400$ GeV, $A_0 = 0$,
  $\tan\beta = 30$, $\mu > 0$.

Global fits in the MSSM

[de Boer, Dabelstein, WH, Mösle, Schwickerath]

[de Boer, Sander]



The Higgs sector of the MSSM

- Two $SU(2) \times U(1)$ doublets: $H_1 = \begin{pmatrix} H_1^0 \\ H_1^- \end{pmatrix}, \quad H_2 = \begin{pmatrix} H_2^+ \\ H_2^0 \end{pmatrix}$

$$H_i^0 = \frac{v_i + S_i + i\,P_i}{\sqrt{2}} \qquad\qquad \tan\beta = \frac{v_2}{v_1}$$

- The soft SUSY-breaking mass terms for H_1^0 and H_2^0 are responsible for electroweak symmetry breaking (EWSB):

$$\begin{aligned} V_{\text{tree}} &= (m_{H_1}^2 + \mu^2)\,|H_1^0|^2 + (m_{H_2}^2 + \mu^2)\,|H_2^0|^2 \\ &+ B\,\left(H_1^0 H_2^0 + \text{h.c.}\right) + \frac{1}{8}(g^2 + g'^2)\,\left(|H_1^0|^2 - |H_2^0|^2\right)^2 \end{aligned}$$

- Five physical states: $h\,,\ H\,,\ A^0\,,\ H^+\,,\ H^-$

- Tree–level mass matrix for the CP–even sector:

$$\left(\mathcal{M}_S^2\right)^{\text{tree}} = \begin{pmatrix} m_Z^2\, c_\beta^2 + m_A^2\, s_\beta^2 & -\left(m_Z^2 + m_A^2\right) s_\beta\, c_\beta \\ -\left(m_Z^2 + m_A^2\right) s_\beta\, c_\beta & m_Z^2\, s_\beta^2 + m_A^2\, c_\beta^2 \end{pmatrix}$$

→ m_h and m_H are predicted in terms of m_Z, m_A and $\tan\beta$

- Tree–level mass relation: $m_h^2 \leq \cos^2 2\beta\, m_Z^2$

dressed Higgs propagators

$$(\Delta_{\text{Higgs}})^{-1} = \begin{pmatrix} q^2 - m_H^2 + \hat{\Sigma}_H(q^2) & \hat{\Sigma}_{hH}(q^2) \\ \hat{\Sigma}_{Hh}(q^2) & q^2 - m_h^2 + \hat{\Sigma}_h(q^2) \end{pmatrix}$$

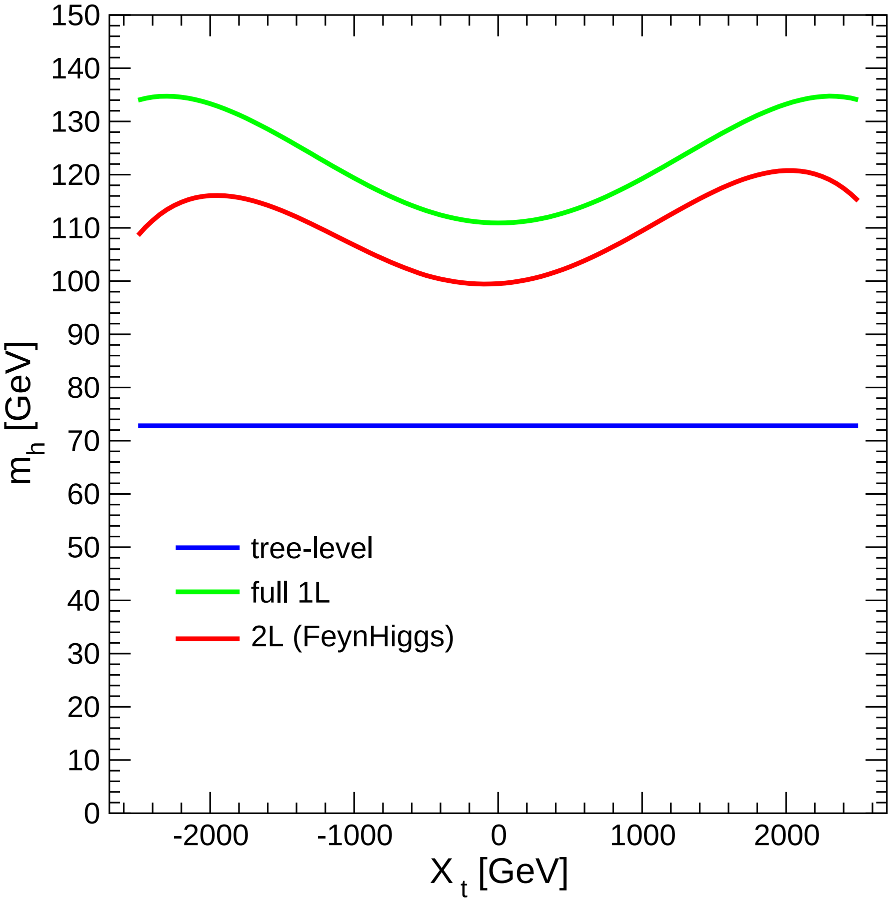
- det = 0 → $m_{h,H}^{\text{pole}}$
- diagonalization → effective couplings (α_{eff})

renormalized self-energies $\hat{\Sigma}$

1-loop: complete

2-loop: QCD corrections $\sim \alpha_s\alpha_t,\ \alpha_s\alpha_b$
Yukawa corrections $\sim \alpha_t^2$
[→ FeynHiggs]

m_{h^0} prediction at different levels of accuracy:



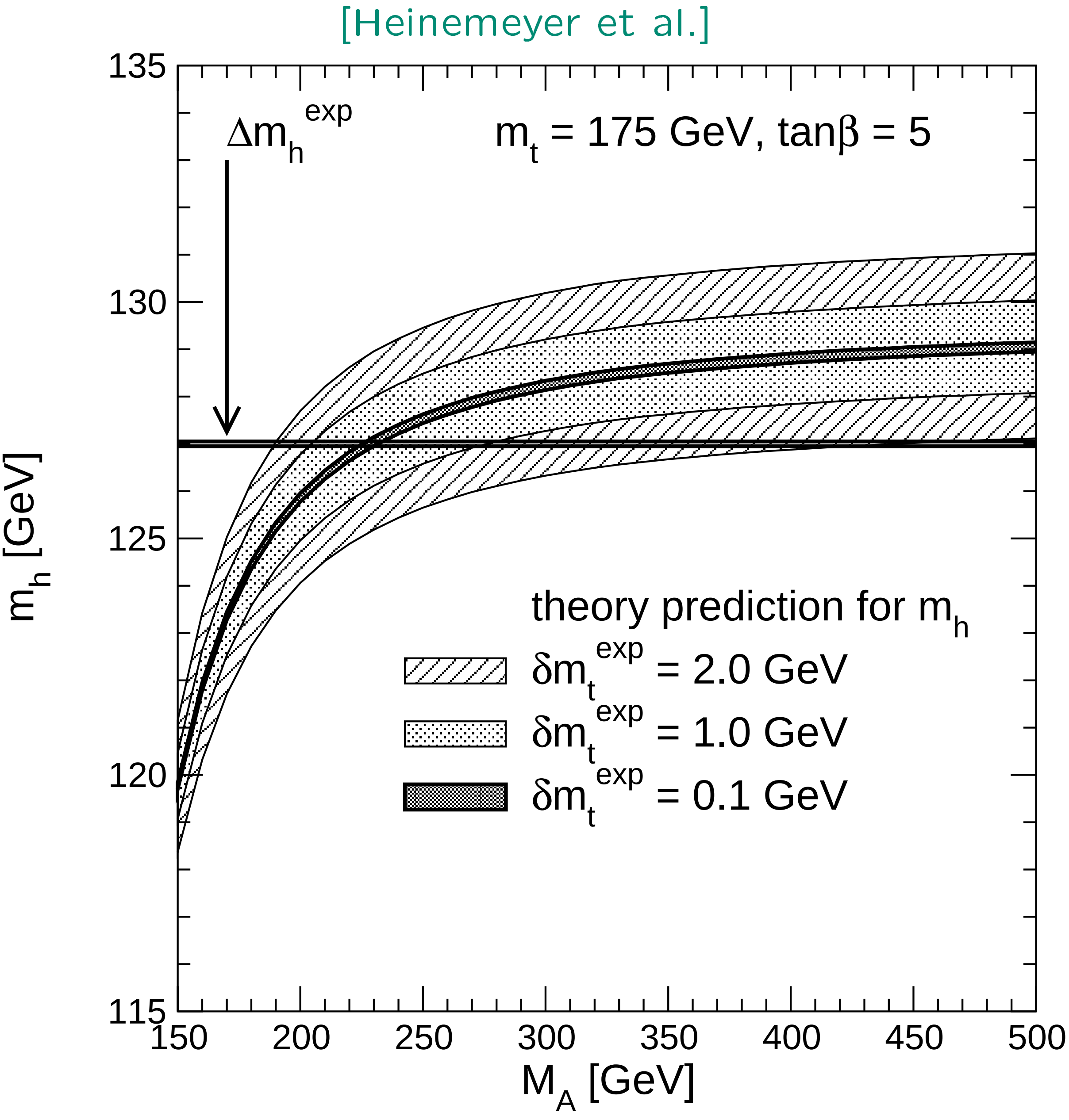
$\tan\beta = 3$, $M_{\tilde{Q}} = M_A = 1$ TeV, $m_{\tilde{g}} = 800$ GeV

X_t : top-squark mixing parameter

$$X_t = A_t - \mu \cot\beta\,, \qquad \mathcal{M}^2_{\tilde{t}} = \begin{pmatrix} m^2_{\tilde{t}_L} & m_t\, X_t \\ m_t\, X_t & m^2_{\tilde{t}_R} \end{pmatrix}$$

present theoretical uncertainty: $\delta m_h \simeq 4$ GeV

[Degrassi, Heinemeyer, WH, Slavich, Weiglein]


[Heinemeyer et al.]
Δm_h^{exp}
m_t = 175 GeV, tanβ = 5
theory prediction for m_h
δm_t^{exp} = 2.0 GeV
δm_t^{exp} = 1.0 GeV
δm_t^{exp} = 0.1 GeV
m_h [GeV]
M_A [GeV]
115
120
125
130
135
150
200
250
300
350
400
450
500

SUSY particles

- LHC will see SUSY if at low energy scale
- LC and LHC$\oplus$LC for precision studies
- Reconstruction of fundamental SUSY theory and breaking mechanism

from experiment:

→ precision analyses of masses and couplings including higher orders

from theory:

→ accurate theoretical predictions to match exp. data

→ loop contributions Lagrangian param ↔ observables

→ RGEs for extrapolation to high scales

chargino/neutralino sector

complete at one loop [Fritzsche,WH/Eberl,Majerotto,...]
renormalization and mass spectrum
pair production in e^+e^- collisions

sfermion sector

renormalization and mass spectrum
[WH, Rzehak]

$$\begin{pmatrix} m_f^2 + M_L^2 + M_Z^2 c_{2\beta}(I_f^3 - Q_f s_W^2) & m_f(A_f - \mu\kappa) \\ m_f(A_f - \mu\kappa) & m_f^2 + M_{\tilde{f}_R}^2 + M_Z^2 c_{2\beta} Q_f s_W^2 \end{pmatrix}$$

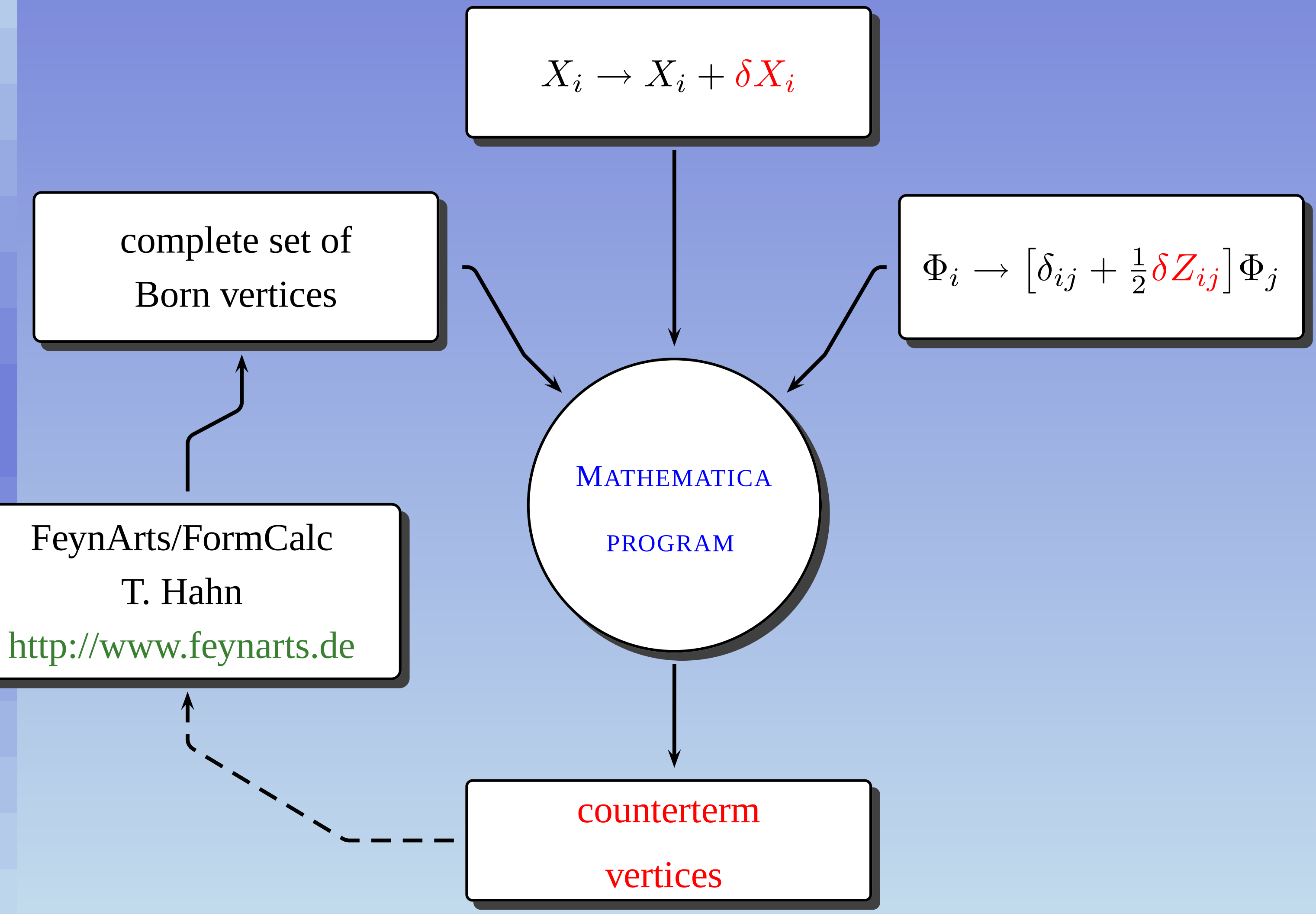
sfermion pair production in e^+e^- collisions
complete at one-loop

[Arhrib, WH]	squarks, sleptons
[Kovarik, Weber, Eberl, Majerotto]	squarks
[Freitas, Miller, von Manteuffel, Zerwas]	sleptons

sfermion decays into fermions and -inos

complete at one-loop
[Guasch, WH, Solà]

Automatic generation of CTs



$$X_i \to X_i + \delta X_i$$

complete set of
Born vertices

$$\Phi_i \to \left[\delta_{ij} + \tfrac{1}{2}\delta Z_{ij}\right]\Phi_j$$

MATHEMATICA
PROGRAM

FeynArts/FormCalc
T. Hahn
http://www.feynarts.de

counterterm
vertices

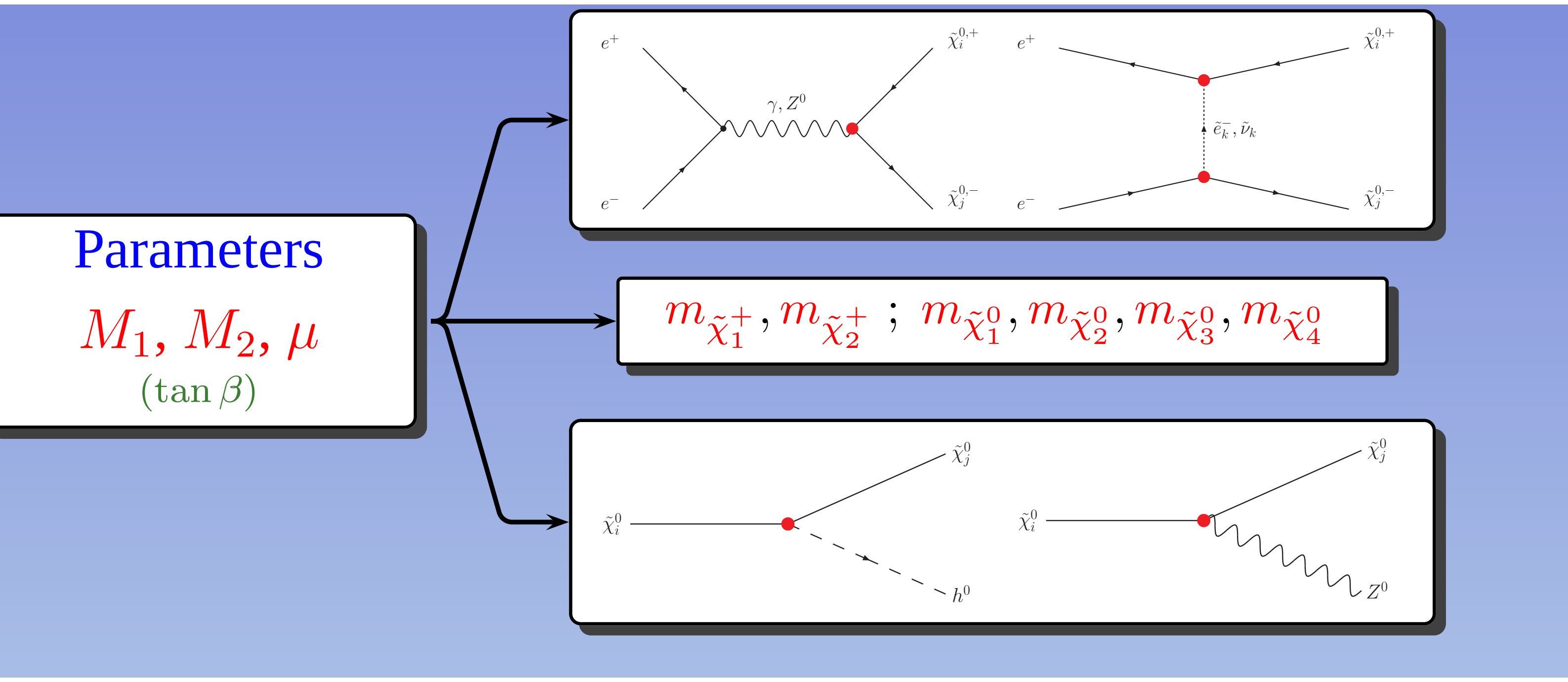
`[FFV]` 2 Charginos – Gauge Boson

$$\underset{30}{C}(-\tilde{\chi}^-_{c2},\tilde{\chi}^-_{c1},\gamma)=\begin{bmatrix} \dfrac{\mathrm{i}\,e\left\{\begin{array}{l} 2\left\{\begin{array}{l} c_W\,s_W\left[\delta Z^{R}_{\tilde{\chi}^+}\right]^{*}_{c1,c2}+c_W\,s_W\left[\delta Z^{R}_{\tilde{\chi}^+}\right]_{c2,c1}+\left(-[\delta Z_{ZA}]_{12}\,s_W{}^2+(c_W\,s_W)\,([\delta Z_{ZA}]_{22}+2\,\delta Z_e)\right)\,\delta_{c1,c2}+\\ [\delta Z_{ZA}]_{12}\,U^*_{c1,1}\,U_{c2,1}\end{array}\right\}+\\ [\delta Z_{ZA}]_{12}\,U^*_{c1,2}\,U_{c2,2}\end{array}\right\}}{4\,c_W\,s_W}\\ \dfrac{\mathrm{i}\,e\left\{\begin{array}{l} 2\left\{\begin{array}{l} c_W\,s_W\left[\delta Z^{L}_{\tilde{\chi}^+}\right]^{*}_{c2,c1}+c_W\,s_W\left[\delta Z^{L}_{\tilde{\chi}^+}\right]_{c1,c2}+\left(-[\delta Z_{ZA}]_{12}\,s_W{}^2+(c_W\,s_W)\,([\delta Z_{ZA}]_{22}+2\,\delta Z_e)\right)\,\delta_{c1,c2}+\\ [\delta Z_{ZA}]_{12}\,V^*_{c2,1}\,V_{c1,1}\end{array}\right\}+\\ [\delta Z_{ZA}]_{12}\,V^*_{c2,2}\,V_{c1,2}\end{array}\right\}}{4\,c_W\,s_W}\end{bmatrix}$$

$$\underset{31}{C}(-\tilde{\chi}^-_{c2},\tilde{\chi}^-_{c1},Z)=\begin{bmatrix} \dfrac{\left\{\begin{array}{l} \left(2\,\delta s_W\,s_W{}^2+c_W{}^2\,(2\,\delta s_W+c_W\,[\delta Z_{ZA}]_{21}-s_W\,([\delta Z_{ZA}]_{11}+2\,\delta Z_e))\right)\,\left(2\,\mathrm{i}\,e\,s_W{}^2\,\delta_{c1,c2}\right)+\\ 4\,c_W{}^3\,s_W{}^2\displaystyle\sum_{n=1}^{2}-\frac{\mathrm{i}\,e\left[\delta Z^{R}_{\tilde{\chi}^+}\right]_{n,c1}}{4\,c_W\,s_W}\,\left(2\,s_W{}^2\,\delta_{n,c2}-2\,U^*_{n,1}\,U_{c2,1}-U^*_{n,2}\,U_{c2,2}\right)+\\ 4\,c_W{}^3\,s_W{}^2\displaystyle\sum_{n=1}^{2}-\frac{\mathrm{i}\,e\left[\delta Z^{R}_{\tilde{\chi}^+}\right]^{*}_{n,c2}}{4\,c_W\,s_W}\,\left(2\,s_W{}^2\,\delta_{c1,n}-2\,U^*_{c1,1}\,U_{n,1}-U^*_{c1,2}\,U_{n,2}\right)+\\ (\mathrm{i}\,e)\,\left(\left(2\,c_W{}^2\,\delta s_W+2\,c_W{}^2\,\delta Z_e\,s_W+c_W{}^2\,[\delta Z_{ZA}]_{11}\,s_W-2\,\delta s_W\,s_W{}^2\right)\,(2\,U^*_{c1,1}\,U_{c2,1}+U^*_{c1,2}\,U_{c2,2})\right)\end{array}\right\}}{4\,c_W{}^3\,s_W{}^2}\\ \dfrac{\left\{\begin{array}{l} \left(2\,\delta s_W\,s_W{}^2+c_W{}^2\,(2\,\delta s_W+c_W\,[\delta Z_{ZA}]_{21}-s_W\,([\delta Z_{ZA}]_{11}+2\,\delta Z_e))\right)\,\left(2\,\mathrm{i}\,e\,s_W{}^2\,\delta_{c1,c2}\right)+\\ 4\,c_W{}^3\,s_W{}^2\displaystyle\sum_{n=1}^{2}-\frac{\mathrm{i}\,e\left[\delta Z^{L}_{\tilde{\chi}^+}\right]_{n,c2}}{4\,c_W\,s_W}\,\left(2\,s_W{}^2\,\delta_{c1,n}-2\,V^*_{n,1}\,V_{c1,1}-V^*_{n,2}\,V_{c1,2}\right)+\\ 4\,c_W{}^3\,s_W{}^2\displaystyle\sum_{n=1}^{2}-\frac{\mathrm{i}\,e\left[\delta Z^{L}_{\tilde{\chi}^+}\right]^{*}_{n,c1}}{4\,c_W\,s_W}\,\left(2\,s_W{}^2\,\delta_{n,c2}-2\,V^*_{c2,1}\,V_{n,1}-V^*_{c2,2}\,V_{n,2}\right)+\\ (\mathrm{i}\,e)\,\left(\left(2\,c_W{}^2\,\delta s_W+2\,c_W{}^2\,\delta Z_e\,s_W+c_W{}^2\,[\delta Z_{ZA}]_{11}\,s_W-2\,\delta s_W\,s_W{}^2\right)\,(2\,V^*_{c2,1}\,V_{c1,1}+V^*_{c2,2}\,V_{c1,2})\right)\end{array}\right\}}{4\,c_W{}^3\,s_W{}^2}\end{bmatrix}$$

`[FFV]` 2 Leptons – Gauge Boson

$$\underset{9}{C}(-e_{j2},e_{j1},\gamma)=\begin{bmatrix}\left([\delta Z_{ZA}]_{12}-2\,[\delta Z_{ZA}]_{12}\,s_W{}^2+(2\,c_W\,s_W)\,([\delta Z_{ZA}]_{22}+2\,\delta Z_e)+4\,c_W\,s_W\,\mathrm{Re}\left[\left[\delta Z^{L}_{e}\right]_{j2,j2}\right]\right)\,\frac{\mathrm{i}\,e\,\delta_{j1,j2}}{4\,c_W\,s_W}\\ \left(-[\delta Z_{ZA}]_{12}\,s_W+c_W\,([\delta Z_{ZA}]_{22}+2\,\delta Z_e)+2\,c_W\,\mathrm{Re}\left[\left[\delta Z^{R}_{e}\right]_{j2,j2}\right]\right)\,\frac{\mathrm{i}\,e\,\delta_{j1,j2}}{2\,c_W}\end{bmatrix}$$



Parameters

M_1, M_2, μ

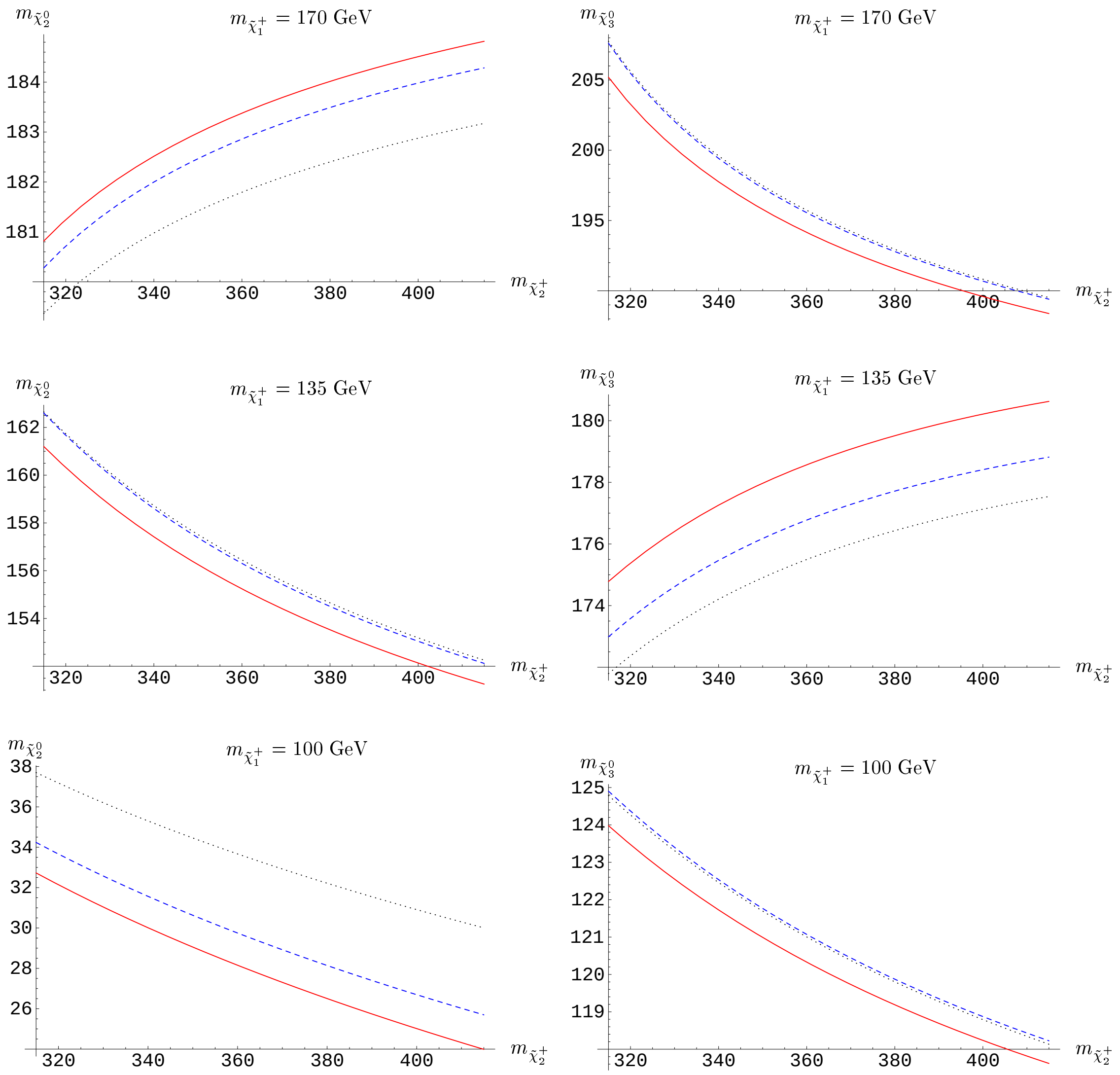
($\tan\beta$)

$m_{\tilde{\chi}_1^+}, m_{\tilde{\chi}_2^+}$; $m_{\tilde{\chi}_1^0}, m_{\tilde{\chi}_2^0}, m_{\tilde{\chi}_3^0}, m_{\tilde{\chi}_4^0}$

pole masses

$[M_{\tilde f} = 300\,\text{GeV},\ \tan\beta = 10]$

$M_{\chi_1^0} = 110\,\text{GeV}$



Born

complete 1-loop

pole masses ↔ on-shell MSSM input

$$\begin{aligned} M_2^2 + \mu^2 + 2M_W^2 &= m_{\tilde{\chi}_1^+}^2 + m_{\tilde{\chi}_2^+}^2 \\ \left(M_2\,\mu - 2M_W^2 \sin\beta\, \cos\beta\right)^2 &= m_{\tilde{\chi}_1^+}^2\, m_{\tilde{\chi}_2^+}^2 \quad . \end{aligned}$$

$$\begin{aligned} M_1 \;=\; & \Big[-M_2 \mu M_Z^2 \sin 2\beta + \big[\mu M_Z^2 \sin 2\beta - M_2\big(\mu^2 + M_Z^2 s_W^2\big)\big] m_{\tilde{\chi}_1^0} \\ & \quad + \big[\mu^2 + M_Z^2\big] m_{\tilde{\chi}_1^0}^2 + M_2 m_{\tilde{\chi}_1^0}^3 - m_{\tilde{\chi}_1^0}^4\Big] \\ & \times \Big[\mu M_Z^2 c_W^2 \sin 2\beta - M_2 \mu^2 + \big[\mu^2 + M_Z^2 c_W^2\big] m_{\tilde{\chi}_1^0} + M_2 m_{\tilde{\chi}_1^0}^2 - m_{\tilde{\chi}_1^0}^3\Big]^{-1} \end{aligned}$$

Renormalization schemes

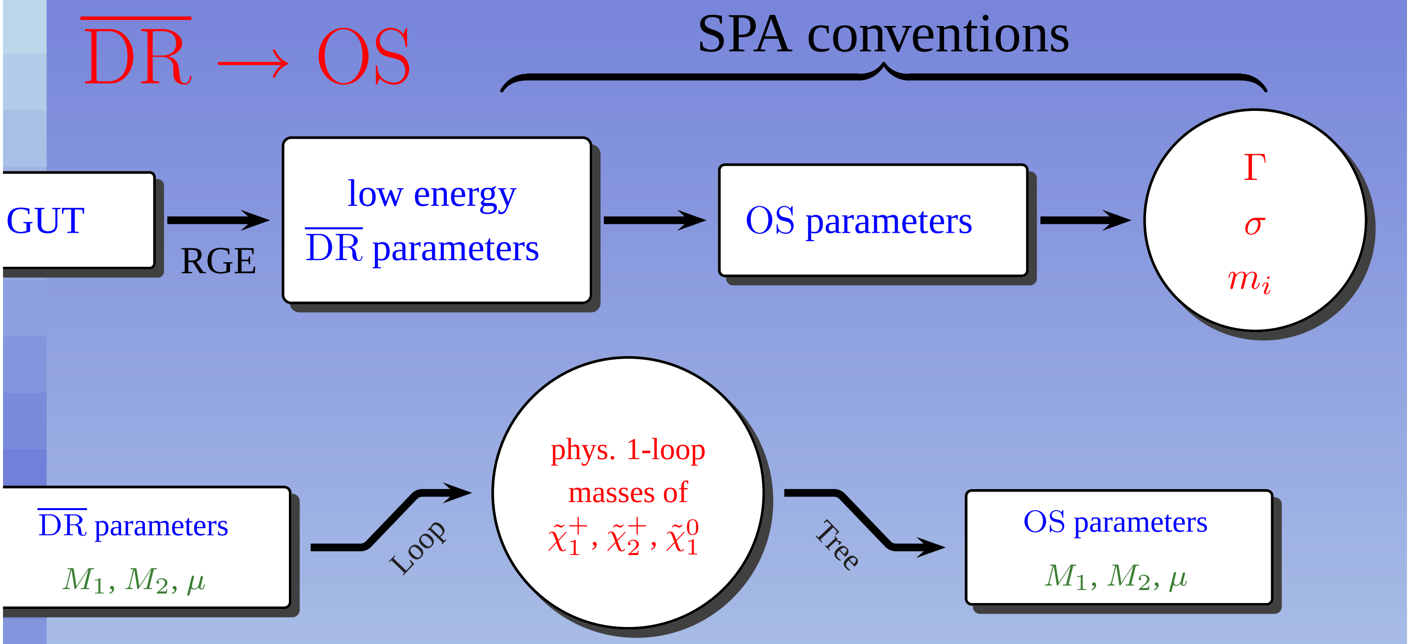
$\overline{\mathrm{DR}}$ scheme:

- Loop integrals: $\frac{2}{\epsilon} - \gamma + \log 4\pi + \log \mu^2 \rightarrow \log \mu^2_{\overline{\mathrm{DR}}}$
- \+ easy to implement
- – observables are scale dependent in finite order perturbation theory
- \+ natural choice for GUT-inspired parameter sets (mSUGRA)

OS scheme:

- renormalization constants fixed by physical conditions
- – renormalization constants complicated
- \+ observables are scale independent
- \+ well suited for calculations of cross sections and decay rates
  (e.g. pole masses → correct kinematical thresholds)

$\overline{\text{DR}} \rightarrow \text{OS}$



SPA conventions

GUT

RGE

low energy $\overline{\text{DR}}$ parameters

OS parameters

Γ
σ
m_i

$\overline{\text{DR}}$ parameters
M_1, M_2, μ

Loop

phys. 1-loop masses of $\tilde{\chi}_1^+, \tilde{\chi}_2^+, \tilde{\chi}_1^0$

Tree

OS parameters
M_1, M_2, μ

```
SPS1a
M1  =  99.1
M2  = 192.7
MUE = 352.4
```

```
MCha(1) = 176.013 +  8.889
MCha(2) = 378.527 + 10.312
MNeu(1) =  96.154 +  4.004
```

```
SPS1a-OS
M1  = 103.02
M2  = 201.56
MUE = 363.06
```

SPA

SUPERSYMMETRY PARAMETER ANALYSIS

The SPA project is a joint study of theorists and experimentalists working on LHC and Linear Collider phenomenology. The study focuses on the supersymmetric extension of the Standard Model. The main targets are

- High-precision determination of the supersymmetry Lagrange parameters at the electroweak scale
- Extrapolation to a high scale to reconstruct the fundamental parameters and the mechanism for supersymmetry breaking

The SPA convention and the SPA Project are described in the report SPA.draft.ps.

`http://spa.desy.de/spa`

P. Zerwas, J. Kalinowski, H.U. Martyn,
W. Hollik, W. Kilian, W. Majerotto,
W. Porod, . . .

SPA CONVENTION

- The masses of the SUSY particles and Higgs bosons are defined as pole masses.
- All SUSY Lagrangian parameters, mass parameters and couplings, including $\tan\beta$, are given in the $\overline{DR}$ scheme and defined at the scale $\tilde{M} = 1$ TeV.
- Gaugino/higgsino and scalar mass matrices, rotation matrices and the corresponding angles are defined in the $\overline{DR}$ scheme at $\tilde{M}$, except for the Higgs system in which the mixing matrix is defined in the on-shell scheme, the scale parameter chosen as the light Higgs mass.
- The Standard Model input parameters of the gauge sector are chosen as G_F, α, M_Z and $\alpha_s^{\overline{MS}}(M_Z)$. All lepton masses are defined on-shell. The t quark mass is defined on-shell; the b, c quark masses are introduced in $\overline{MS}$ at the scale of the masses themselves while taken at a renormalization scale of 2 GeV for the light u, d, s quarks.
- Decay widths / branching ratios and production cross sections are calculated for the set of parameters specified above.

Renormalization of $\tan\beta$

from the Higgs sector

$$\tan\beta = \frac{v_2}{v_1} \;\rightarrow\; \sqrt{\frac{Z_{H_2}}{Z_{H_1}}} \cdot \frac{v_2 + \delta v_2}{v_1 + \delta v_1}$$

$$= \frac{v_2}{v_1}\left(1 + \underset{\overline{DR}}{\delta Z_{H_2} - \delta Z_{H_1}} + \underset{= 0}{\frac{\delta v_2}{v_2} - \frac{\delta v_1}{v_1}}\right)$$

Separation of "QED corrections"

Full calculation inevitable

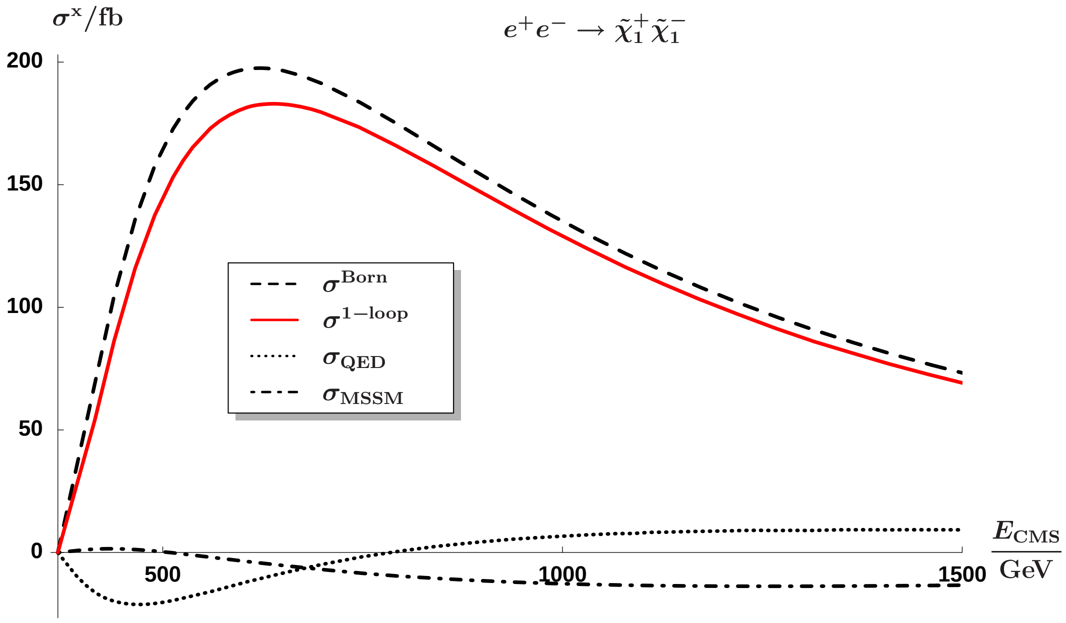
- separation of diagrams with virtual photons not UV-finite
- soft-photon bremsstrahlung necessary for IR-finite result
- hard bremsstrahlung for realistic treatments

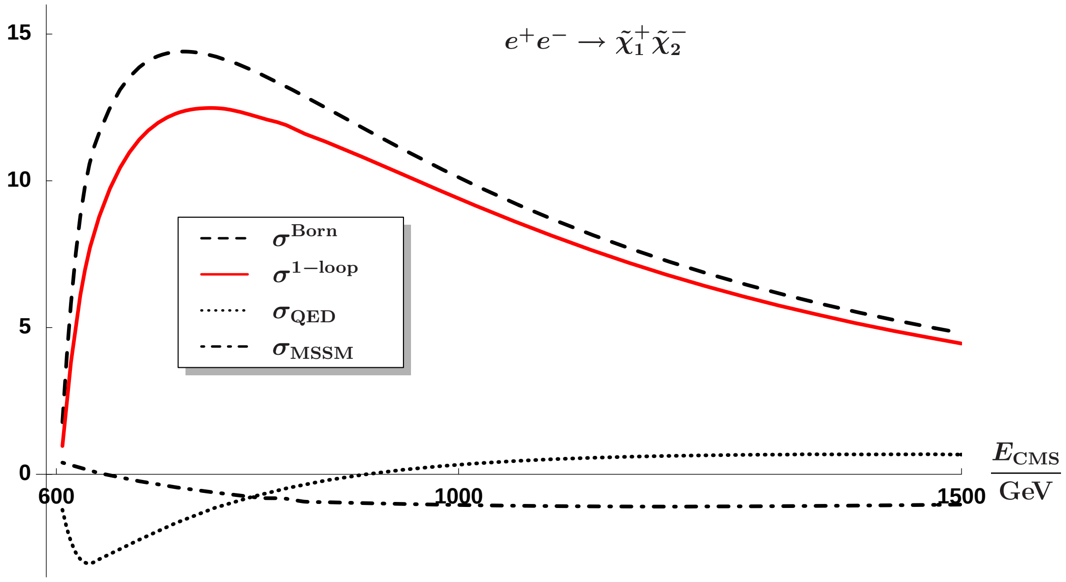
Reasonable separation $(L_e = \log \frac{s}{m_e^2},\ \Delta E = E_{\gamma\,\mathrm{soft}}^{\mathrm{max}})$:

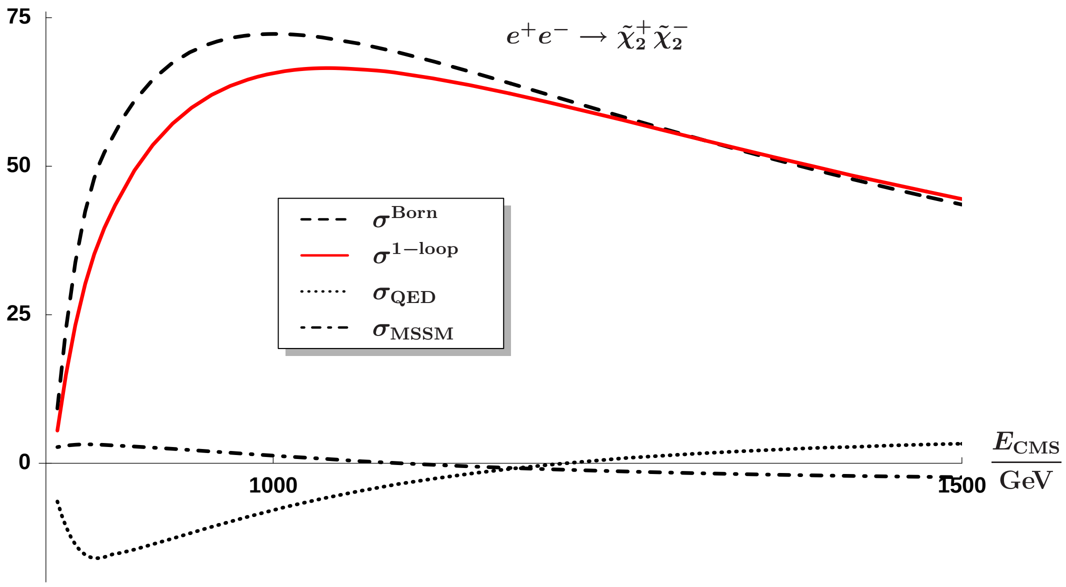
$$\sigma = \sigma_{\mathrm{QED}} + \sigma_{\mathrm{MSSM}}\,,$$

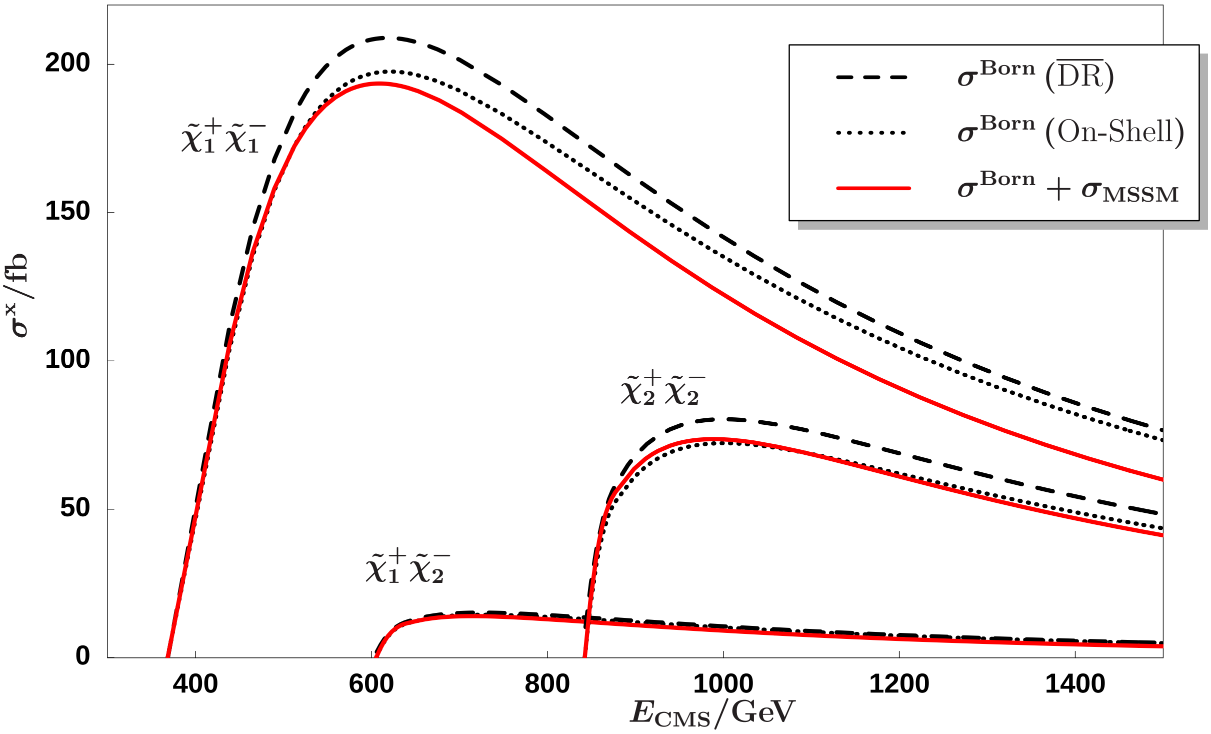
$$\begin{aligned}
\sigma_{\mathrm{QED}} &= \sigma^{\mathrm{hard}} + \frac{\alpha}{\pi}\left[(L_e - 1)\log\frac{4\,\Delta E^{\,2}}{s} + \frac{3}{2}\,L_e\right]\,\sigma_0 \\
\sigma_{\mathrm{MSSM}} &= \sigma^{\mathrm{v+s}} - \frac{\alpha}{\pi}\left[(L_e - 1)\log\frac{4\,\Delta E^{\,2}}{s} + \frac{3}{2}\,L_e\right]\,\sigma_0
\end{aligned}$$

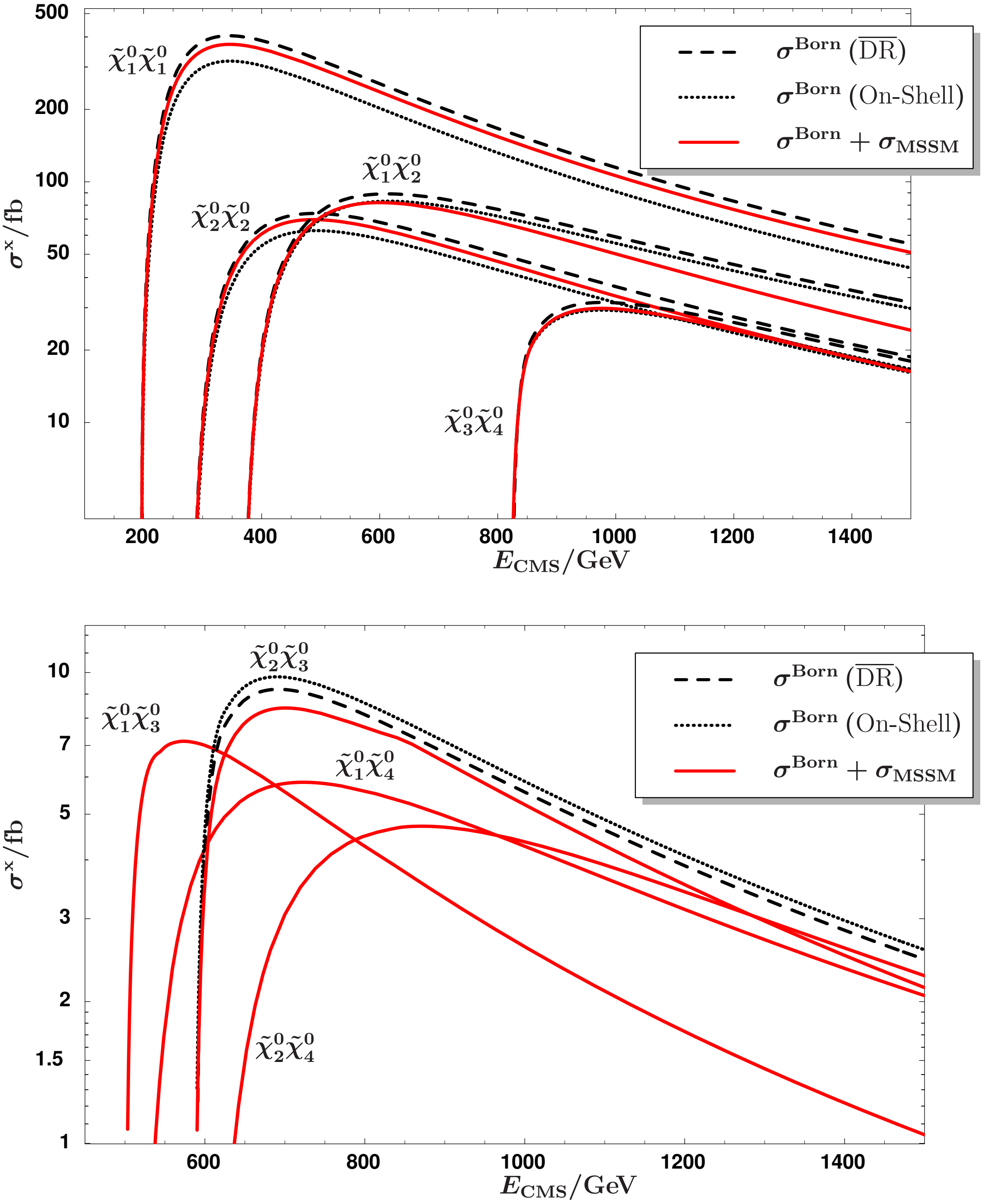
- gauge invariant
- σ_{MSSM} free of large soft and collinear photon contributions


σ^{x}/fb
$e^+e^- \to \tilde{\chi}_1^+\tilde{\chi}_1^-$
200
150
100
50
0
500
1000
1500
$\frac{E_{\mathrm{CMS}}}{\mathrm{GeV}}$
σ^{Born}
$\sigma^{1-\mathrm{loop}}$
σ_{QED}
σ_{MSSM}


$e^+e^- \to \tilde{\chi}_1^+\tilde{\chi}_2^-$
15
10
5
0
600
1000
1500
$\frac{E_{\mathrm{CMS}}}{\mathrm{GeV}}$
σ^{Born}
$\sigma^{1-\mathrm{loop}}$
σ_{QED}
σ_{MSSM}


$e^+e^- \to \tilde{\chi}_2^+\tilde{\chi}_2^-$
75
50
25
0
1000
1500
$\frac{E_{\mathrm{CMS}}}{\mathrm{GeV}}$
σ^{Born}
$\sigma^{1-\mathrm{loop}}$
σ_{QED}
σ_{MSSM}


σ^{Born} ($\overline{\mathrm{DR}}$)
σ^{Born} (On-Shell)
$\sigma^{\mathrm{Born}} + \sigma_{\mathrm{MSSM}}$
$\tilde{\chi}_1^+\tilde{\chi}_1^-$
$\tilde{\chi}_2^+\tilde{\chi}_2^-$
$\tilde{\chi}_1^+\tilde{\chi}_2^-$
σ^{x}/fb
0
50
100
150
200
400
600
800
1000
1200
1400
$E_{\mathrm{CMS}}/\mathrm{GeV}$



$e^+e^- \to \tilde{\chi}_i^0\tilde{\chi}_j^0$

REFERENCE POINT SPS1a′

SPS1a′ deriv. of Snowmass Point SPS1a: conform with Ω_{cdm}, LE data

mSUGRA values:

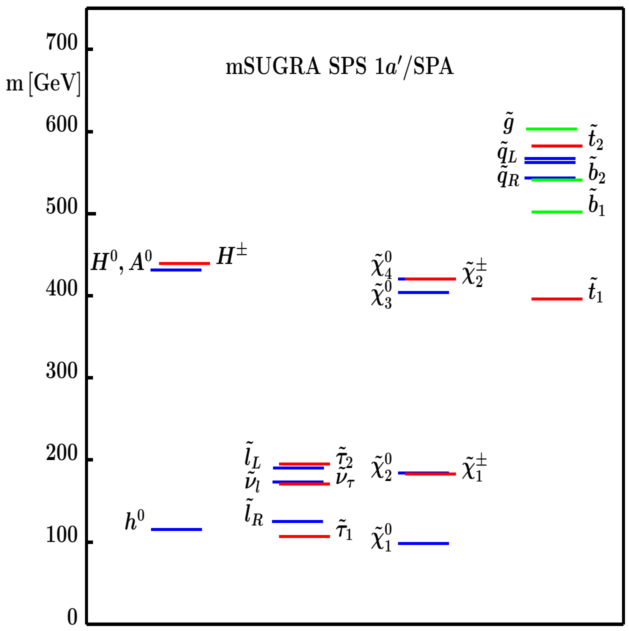
$$\begin{array}{lcr lcr} M_{1/2} & = & 250\text{ GeV} & \text{sign}(\mu) & = & +1 \\ M_0 & = & 70\text{ GeV} & \tan\beta & = & 10 \\ A_0 & = & -300\text{ GeV} & & & \end{array}$$

LE/cosmic parameters:

$BR(b \to s\gamma) = 3.0 \times 10^{-4}$	micrOMEGAs
$\Delta[g_\mu - 2]/2 = 34 \times 10^{-10}$	FeynHiggs
$\Omega_{cdm} h^2 = 0.10$	micrOMEGAs

POLE MASSES:

	m [GeV]		m [GeV]
h^0	115.4	$\tilde{e}_R$	125.2
H^0	431.1	$\tilde{e}_L$	190.1
A^0	431.0	$\tilde{\nu}_e$	172.8
H^+	438.6	$\tilde{\tau}_1$	107.4
$\tilde{\chi}_1^0$	97.75	$\tilde{\tau}_2$	195.3
$\tilde{\chi}_2^0$	184.4	$\tilde{\nu}_\tau$	170.7
$\tilde{\chi}_3^0$	406.8	$\tilde{u}_R$	547.7
$\tilde{\chi}_4^0$	419.6	$\tilde{u}_L$	565.7
$\tilde{\chi}_1^+$	184.2	$\tilde{t}_1$	368.9
$\tilde{\chi}_2^+$	421.1	$\tilde{t}_2$	584.9
$\tilde{g}$	607.6	$\tilde{b}_1$	506.3



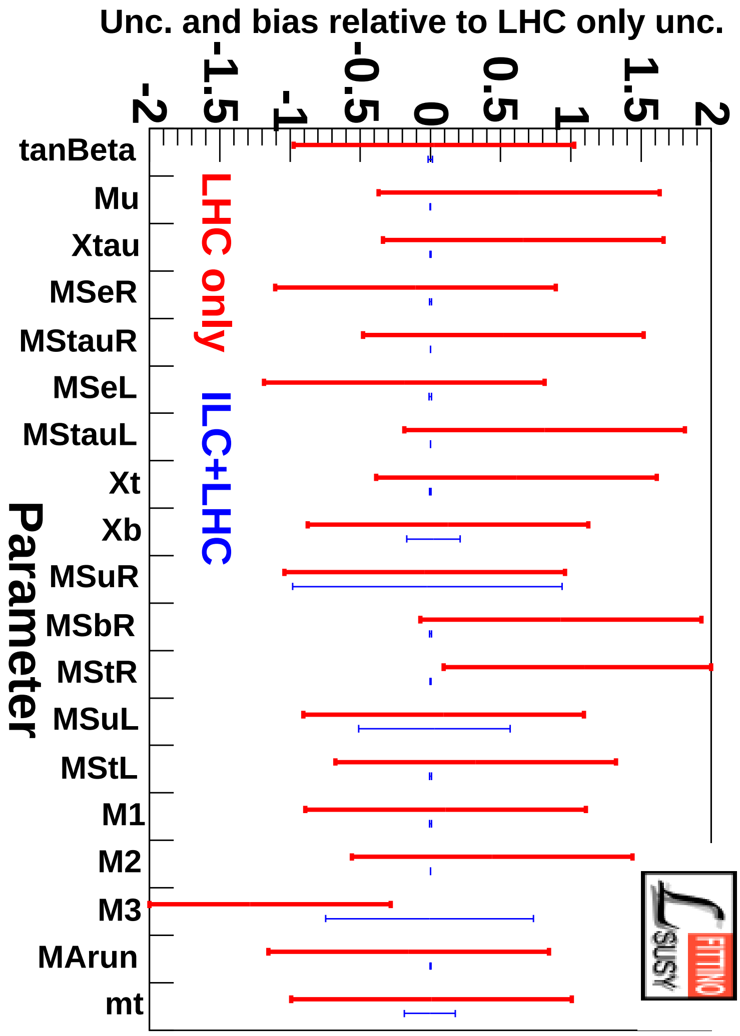
$BR(\tilde{\nu} \to \nu \chi_1^0) = 100\%$ $\Rightarrow$ $\tilde{\nu}$ invis.

Measurements

- edge effects at LHC
- decay spectra at ILC
- cross sections/asymmetries at ILC

	Mass	"LHC"	"LC"	"LHC+LC"
h^0	115.4	0.25	0.05	0.05
H^0	431.1		1.5	1.5
$\tilde{\chi}_1^0$	97.75	4.8	0.05	0.05
$\tilde{\chi}_2^0$	184.4	4.7	1.2	0.08
$\tilde{\chi}_4^0$	419.6	5.1	3 – 5	2.5
$\tilde{\chi}_1^\pm$	184.2		0.55	0.55
$\tilde{e}_R$	125.2	4.8	0.05	0.05
$\tilde{e}_L$	190.1	5.0	0.18	0.18
$\tilde{\tau}_1$	107.4	5 – 8	0.24	0.24
$\tilde{q}_R$	547.7	7 – 12	–	5 – 11
$\tilde{q}_L$	565.7	8.7	–	4.9
$\tilde{t}_1$	368.9		1.9	1.9
$\tilde{b}_1$	506.3	7.5	–	5.7
$\tilde{g}$	607.6	8.0	–	6.5

[Bechtle, LCWS 05]



Reconstructing Lagrange param.

based on 82 simulated measurements at LHC and ILC

Parameter	SPS1a$'$value	Fit error [exp]
M_1	103.3	0.1
M_2	193.4	0.1
M_3	568.9	7.8
μ	400.4	1.1
$M_{\tilde{e}_L}$	181.3	0.2
$M_{\tilde{e}_R}$	115.6	0.4
$M_{\tilde{\tau}_L}$	179.5	1.2
$M_{\tilde{u}_L}$	523.2	5.2
$M_{\tilde{u}_R}$	503.9	17.3
$M_{\tilde{t}_L}$	467.7	4.9
m_{A}	374.9	0.8
A_{t}	-525.6	24.6
$\tan\beta$	10.0	0.3

High Scale Extrapolations

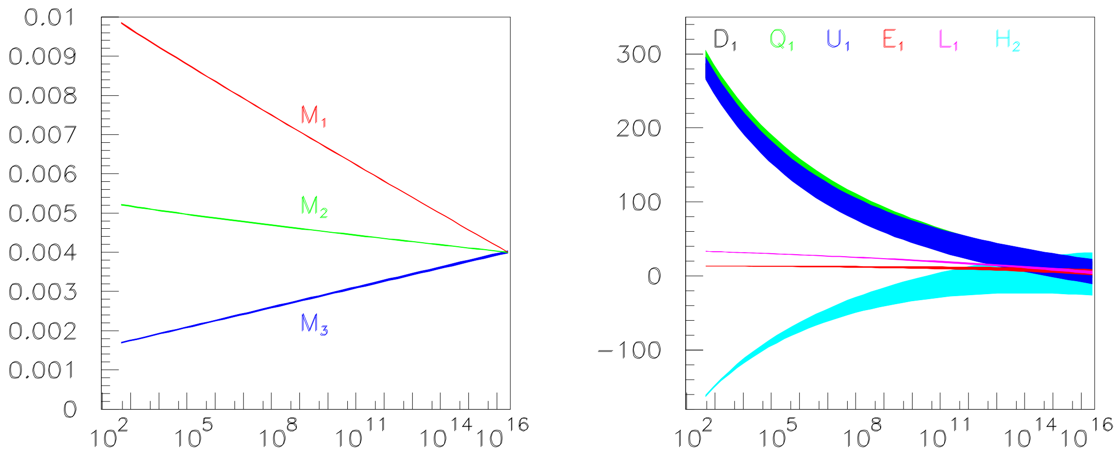


Fig. 1. *Running of the gaugino and scalar mass parameters in SPS1a′ [*`SPheno 2.2.2`*]. Only experimental errors are taken into account; theoretical errors are assumed to be reduced to the same size in the future.*

ERRORS `SPS1a′`:

mSUGRA	Parameter, ideal	"LHC+LC" errors
M_1	250. GeV	0.18 GeV
M_2	*ditto*	0.26 GeV
M_3		2.8 GeV
M_{L_1}	70. GeV	4.1 GeV
M_{E_1}	*ditto*	7.9 GeV
M_{Q_1}		11. GeV
M_{U_1}		31. GeV
M_{H_1}	*ditto*	7.5 GeV
M_{H_2}		72. GeV
A_t	−300. GeV	44. GeV

CONCLUSION:

- gauginos in excellent $\mathcal{O}$[per-mille] condition
- scalar leptons in good $\mathcal{O}$[per-cent] condition
- squarks in $\mathcal{O}[1]$ condition

Conclusions

- Era of electroweak precision physics:
  - quantum effects have been established
  - strong indication for a light Higgs boson

- The MSSM is competitive to the SM
  - global fits of similar quality (slightly better)
  - natural: light Higgs boson

- m_{h^0} is another precision observable
  - dependent on all SUSY sectors
  - accurate theoretical evaluation ($\delta m_{h^0} \simeq 4$ GeV), to be further improved

- one-loop studies for SUSY processes are underway, many results and tools already available