

Prospects for a Gravitino LSP in Gaugino Mediation

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Based on

W. Buchmüller, K. Hamaguchi, J.K., hep-ph/0506105

Outline

- 1 Gaugino Mediation
- 2 Naive Dimensional Analysis
- 3 Bounds on Gaugino and Gravitino Masses

Motivation

Gravitino LSP

- Alternative dark matter candidate
- Can allow large enough reheating temperature for thermal leptogenesis
- Decays of a charged NLSP may be observable at colliders

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Gravitino mass in different SUSY-breaking scenarios

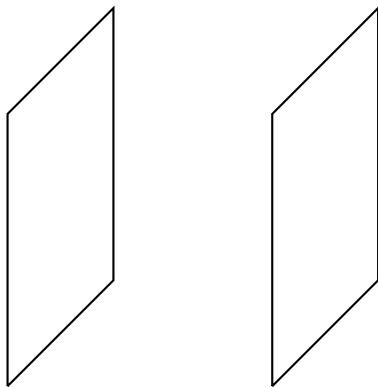
- Gauge mediation: $m_{3/2} \ll 1 \text{ GeV} \Rightarrow \text{LSP}$
- Gravity mediation: $m_{3/2} \sim 100 \text{ GeV} - 1 \text{ TeV} \Rightarrow \text{maybe LSP}$
- Anomaly mediation: $m_{3/2} \gtrsim 10 \text{ TeV} \Rightarrow \text{not LSP}$
- Gaugino mediation: ?

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Setup

Kaplan, Kribs, Schmaltz, Phys. Rev. **D62** (2000)

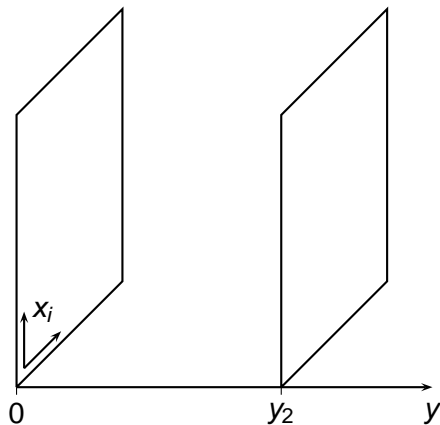
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- D dimensions
- 4-dimensional branes

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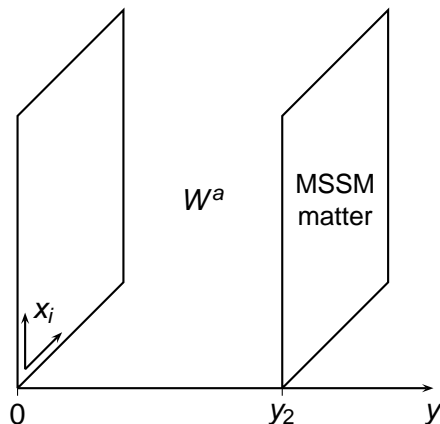
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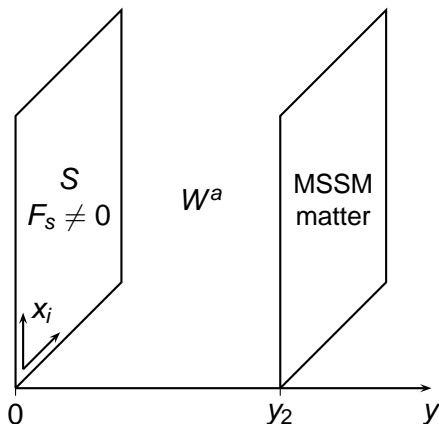
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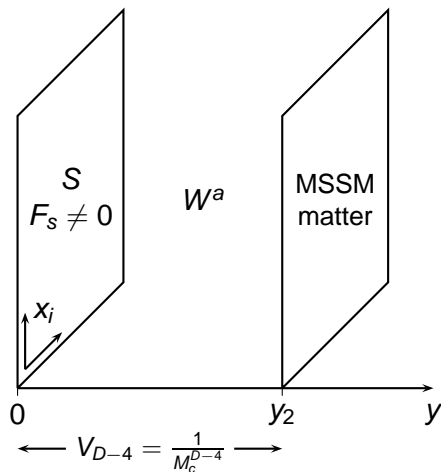
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Coupling of Gauge Fields to S

$$\mathcal{L}_{SWW} = \delta^{(D-4)}(y) \frac{h}{4\Lambda} \int d^2\theta S W^a W^a + \text{h.c.}$$

- h : dimensionless coupling
- Λ : cutoff

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- Λ : cutoff
- Normalization of gauge fields:

$$\mathcal{L}_{\text{kin}} = \frac{1}{4g_D^2} \int d^2\theta W^a W^a + \text{h.c.} + \dots$$

SUSY Breaking

Vev F_S for the F -term of S yields

- Gaugino mass

$$m_{1/2} = \frac{g_4^2 h F_S}{2\Lambda}$$

- g_4 : 4-dimensional gauge coupling ($g_4^2 = \frac{g_D^2}{V_{D-4}}$)
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- Squark and slepton masses very small
 \rightsquigarrow solution of the SUSY flavor problem

Restricting the Masses

$$\frac{m_{3/2}}{m_{1/2}} = \frac{2\Lambda}{\sqrt{3}g_4^2 h M_4}$$

- $M_c < \Lambda \lesssim M_D$
- $g_4^2 \sim \frac{1}{2}$ for $M_c \sim M_{\text{GUT}}$

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 - Must not be too large to avoid strong coupling below Λ
 \Rightarrow **Lower bound** on $m_{3/2}/m_{1/2}$
 \rightsquigarrow Estimate by naive dimensional analysis

- 1 Gaugino Mediation
- 2 Naive Dimensional Analysis**
- 3 Bounds on Gaugino and Gravitino Masses

Idea

Manohar, Georgi, Nucl. Phys. **B234** (1984); Georgi, Randall, Nucl. Phys. **B276** (1986)
 Luty, Phys. Rev. **D57** (1998); Cohen, Kaplan, Nelson, Phys. Lett. **B412** (1997)
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Consider theory with **dimensionless** bulk fields $\hat{\Phi}(x, y)$ and brane fields $\hat{\phi}(x)$, valid up to a cutoff Λ :

$$\mathcal{L} = \Lambda^D \mathcal{L}_{\text{bulk}}(\hat{\Phi}) + \delta^{(D-4)}(y) \Lambda^4 \mathcal{L}_1(\hat{\Phi}, \hat{\phi})$$

Loop corrections

- suppressed by **loop factor** $1/l_D$ ($l_4 = 16\pi^2$, $l_5 = 24\pi^3$, ...)
- enhanced by **group theory factor** C , typically $C \sim C_2(G)$ (e.g. $C = 5$ for $SU(5)$)

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$$\mathcal{L} = \frac{\Lambda^D}{\ell_D/C} \hat{\mathcal{L}}_{\text{bulk}}(\hat{\Phi}) + \delta^{(D-4)}(y) \frac{\Lambda^4}{\ell_4/C} \hat{\mathcal{L}}_1(\hat{\Phi}, \hat{\phi})$$

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⇒ With suitable prefactors in \mathcal{L} :

All loops of same order of magnitude for $\mathcal{O}(1)$ couplings

⇔ **Strong coupling**

NDA Recipe

- 1 Express theory in terms of dimensionless fields:

$$\mathcal{L} = \frac{\Lambda^D}{l_D/\mathbf{C}} \hat{\mathcal{L}}_{\text{bulk}}(\hat{\Phi}) + \delta^{(D-4)}(y) \frac{\Lambda^4}{l_4/\mathbf{C}} \hat{\mathcal{L}}_1(\hat{\Phi}, \hat{\phi})$$

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Example

5D gauge coupling:

$$g_5^2 < \frac{24\pi^3/C}{\Lambda}$$

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Constraining the SWW Coupling

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 &= \delta^{(D-4)}(y) \frac{\Lambda^4}{l_4/C} \frac{g_D^2 h \sqrt{l_4 C} \Lambda^{D-4}}{l_D} \frac{1}{4} \int \frac{d^2\theta}{\Lambda} \hat{S} \hat{W}^a \hat{W}^a + \text{h.c.}
 \end{aligned}$$

Constraining the SWW Coupling

$$\begin{aligned}
 \mathcal{L}_{SWW} &= \delta^{(D-4)}(y) \frac{h}{4\Lambda} \int d^2\theta \, S W^a W^a + \text{h.c.} \\
 &= \delta^{(D-4)}(y) \frac{\Lambda^4}{l_4/C} \underbrace{\frac{g_D^2 h \sqrt{l_4 C} \Lambda^{D-4}}{l_D}}_{< 1} \frac{1}{4} \int \frac{d^2\theta}{\Lambda} \hat{S} \hat{W}^a \hat{W}^a + \text{h.c.} \\
 \Rightarrow h &< \frac{l_D}{4\pi\sqrt{C}g_4^2} \left(\frac{M_C}{\Lambda}\right)^{D-4}
 \end{aligned}$$

Constraining the SWW Coupling

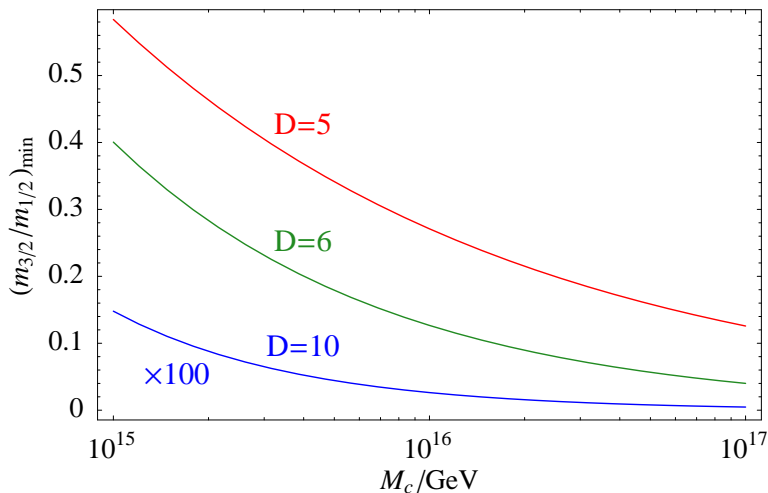
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 &\Rightarrow h < \frac{l_D}{4\pi\sqrt{C}g_4^2} \left(\frac{M_c}{\Lambda}\right)^{D-4} \\
 &\quad \Downarrow
 \end{aligned}$$

Lower Bound on Gravitino to Gaugino Mass Ratio

$$\frac{m_{3/2}}{m_{1/2}} = \frac{2\Lambda}{\sqrt{3}g_4^2 h M_4} > \frac{8\pi\sqrt{C}}{\sqrt{3}l_D} \left(\frac{\Lambda}{M_c}\right)^{D-4} \frac{\Lambda}{M_4}$$

Lower Bound on $m_{3/2}/m_{1/2}$

Choose cutoff equal to D -dimensional Planck scale and $C = 5$



Lower Bound on $m_{3/2}/m_{1/2}$

Result

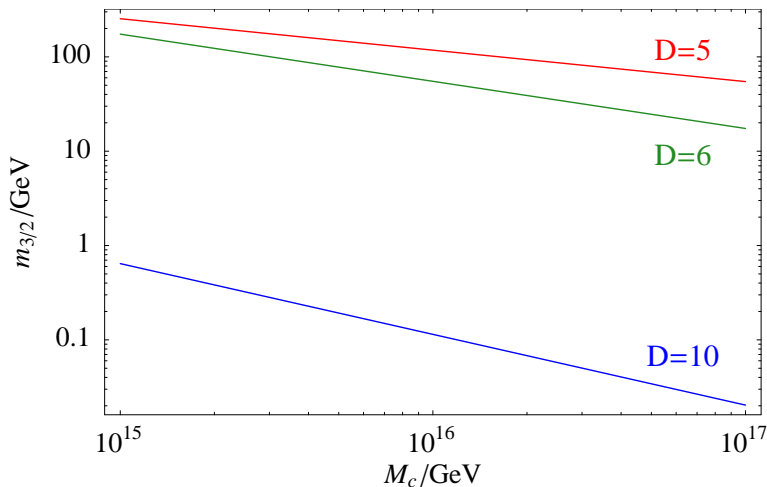
- Gravitino can be the LSP, if gaugino mass sufficiently close to upper bound
- Gravitino can't be much lighter than gauginos for $D = 5, 6$

Note:

- $m_{1/2}$ given at $M_c \rightsquigarrow$ Running to low energies decreases lightest gaugino mass by factor ~ 0.4
- Bound becomes weaker if the cutoff is lowered

Lower Bound on $m_{3/2}$

Fix $m_{\tilde{g}} = 1$ TeV (at low energy) \Rightarrow Lower bound on $m_{3/2}$



Lower Bound on $m_{3/2}$

Typical Bound in Gaugino Mediation

$m_{3/2} > 17 \text{ GeV}$ (for $D = 6$ and $M_c = 10^{17} \text{ GeV}$)

⇒ Experimental implications for much lighter gravitino would disfavor gaugino mediation

Scalar Masses

- Initial condition at compactification scale:
Squark and slepton masses ≈ 0
 - Generated by radiative corrections
 - Lightest slepton is a stau
 - Typical mass slightly below that of lightest neutralino
- ⇒ Gravitino LSP + $\tilde{\tau}$ NLSP possible
- ⇒ $\tilde{\tau}$ decays into gravitinos may be observed at colliders

Buchmüller, Hamaguchi, Ratz, Yanagida, Phys. Lett. **B588** (2004)

Hamaguchi, Kuno, Nakaya, Nojiri, Phys. Rev. **D70** (2004)

Feng, Smith, Phys. Rev. **D71** (2005)

Hamaguchi, Ibarra, JHEP **02** (2005)

Brandenburg, Covi, Hamaguchi, Roszkowski, Steffen, Phys. Lett. **B617** (2005)

Conclusions

For gaugino mediated SUSY breaking:

- Lower bound on mass ratio $m_{3/2}/m_{1/2}$ from naive dimensional analysis
- Gravitino can be the LSP
- Lower bound: $m_{3/2} \gtrsim 10 \text{ GeV}$
- Stau is a good candidate for the NLSP