

# Growth of a Susy Bubble in Dense Matter

(Physics at the high density frontier)

Susy 2005

Durham, July 2005

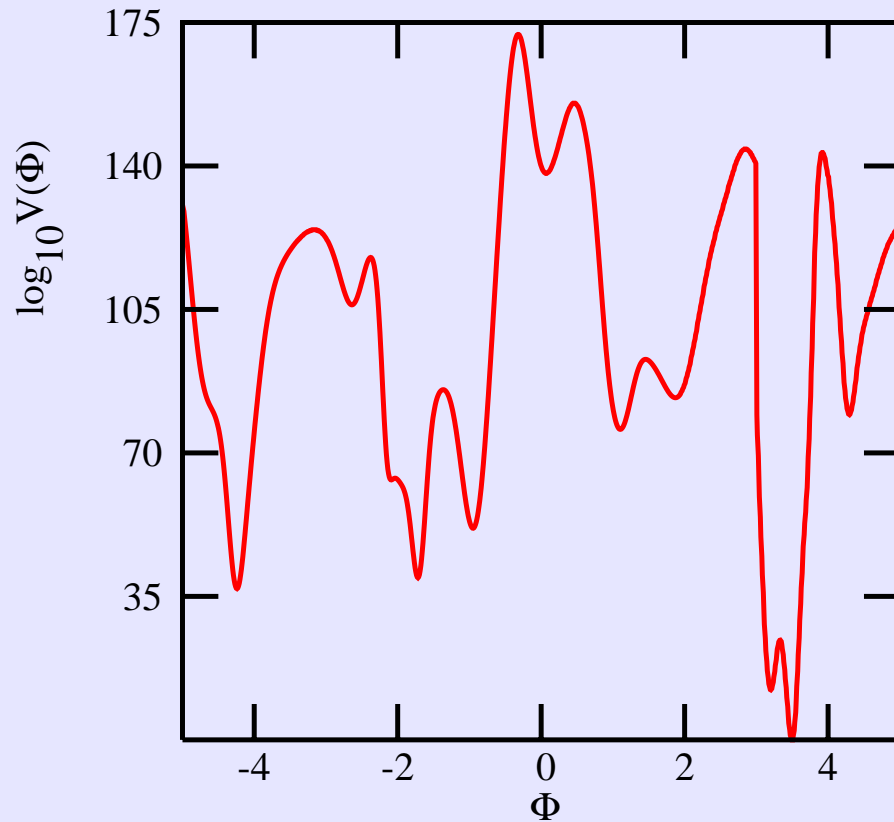
- It is not clear whether the standard model has within it an adequate energy release mechanism to explain supernovae and gamma ray bursts.
- Does susy have a role to play here?
- The matter densities prevailing in white dwarf and neutron stars are millions of times greater than obtainable in bulk matter on earth.

L Clavelli, U Alabama

## Characteristics of Gamma Ray Bursts

- **dominant photon energies:**  
 $\approx 100 \text{ KeV to } 1 \text{ MeV}$
- **total energy in burst:  $\approx 5 \cdot 10^{50}$  ergs**  
assuming  $5^\circ$  jet opening angle  
(earth rest energy =  $6 \cdot 10^{48}$  ergs)  
(sun rest energy =  $2 \cdot 10^{54}$  ergs)
- **total duration:**  
 $0.02s < \tau < 2s$  (short duration bursts)  
 $2s < \tau < 300s$  (long duration bursts)
- **burst rate:**  
 $\frac{dN}{dt} \approx 5 \cdot 10^{-7} \text{ yr}^{-1} \text{ gal}^{-1}$   
(but much less in our galactic neighborhood)

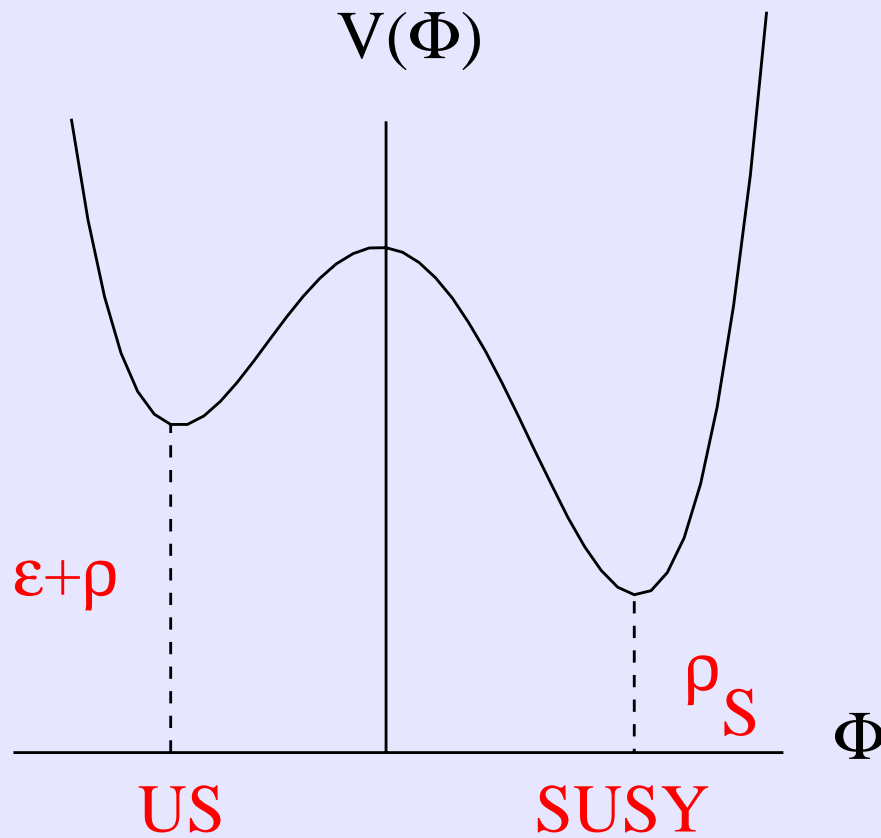
## String Landscape



vacuum energy density:  $\epsilon = 3560 \text{ MeV}/\text{m}^3 = (0.0023 \text{ eV})^4$

Planck density:  $M_{Pl}^4 = 10^{127} \text{ MeV}/\text{m}^3 = (1.66 \cdot 10^{28} \text{ eV})^4$

Effective potential  
in dense media



We live in a broken susy phase with positive vacuum energy  
while the basic string theories predict  
an exact susy ground state with vanishing vacuum energy

## **Gamma Rays from a Susy Star**

**In our broken susy world, the Pauli principle functions as an enormous energy storage mechanism.**

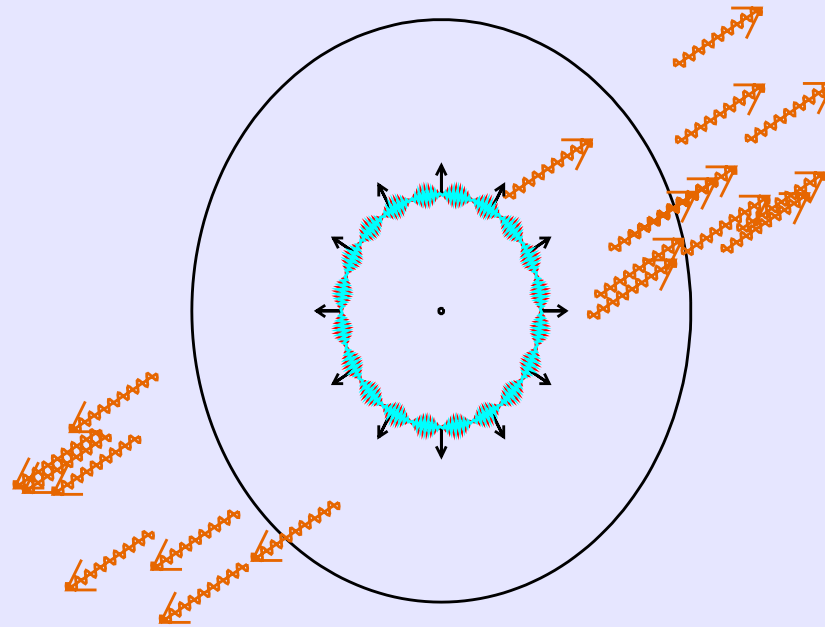
**A phase transition to exact susy would allow the recapture of this energy.**

## Proposed Model for Gamma Ray Bursts

LC+G. Karatheodoris, hep-ph/0403227

LC + I. Perevalova, Phys Rev D, March 05

1. A compact star with a high level of fermion degeneracy (white dwarf or neutron star) makes a phase transition to a state of exact Supersymmetry.
2. Electron pairs undergo quasi elastic scattering to selectron pairs via photino exchange.  $ee \rightarrow \tilde{e}\tilde{e}$
3. Uninhibited by the Pauli principle, the selectrons fall into the ground state emitting photons which can penetrate into the broken-SUSY world.
4. Jet structure is aided by stimulated emission of sfermions and gammas.
5. With no further support from electron degeneracy, the star collapses below the Schwarzschild radius and becomes a black hole.



**a resonant cavity of susy particles  
slowly growing to engulf the star  
emitting gamma rays which are light in both phases**

broken susy phase with positive vacuum energy density:  $\epsilon = 3560 \text{MeV}/m^3$

effective potential for susy bubble in vacuo:  $V(vac) = -4\pi(r^3\epsilon/3 - r^2S)$

$$R_c(vac) = \frac{3S}{\epsilon}$$

*transition prob* :  $\frac{d^2P}{dt d^3r} = A e^{-B(vac)}$

**Exact Susy: the physics of the future**

$$B(vac) = \frac{\pi^2}{6} R_c^4 \epsilon = \frac{27\pi^2 S^4}{2\epsilon^3}$$

$$R_c(vac) > R_{galaxy} \approx 4.7 \cdot 10^{20} \text{ m}$$

$$S > 5.6 \cdot 10^{23} \text{MeV}/m^2 = 2 \cdot 10^{-23} M_\odot R_E^{-2} = (0.28 \text{MeV})^3$$

in dense matter (plausibly):  $\epsilon \rightarrow \epsilon + \rho - \rho_S = \epsilon + \Delta\rho$



$$B = \frac{27\pi^2 S^4}{2(\epsilon + \Delta\rho)^3}$$

$$p_F = \left( \frac{3\pi^2 N}{V} \right)^{1/3}$$

$$N/V = \frac{\rho}{2M_N}$$

$$\langle E \rangle - m = m \left( -1 + {}_2F_1\left(-1/2, 3/2; 5/2; -p_F^2/m^2\right) \right)$$

limit of zero electron mass:  $\langle E \rangle = \frac{3p_F}{4}$

$$\Delta\rho = \frac{3p_F N}{4V} = \frac{1}{4\pi^2} \left( \frac{3\pi^2 \rho}{2M_N} \right)^{4/3}$$

transition probability :  $\frac{d^2 P}{dt d^3 r} = A e^{-B}$

$$B = \left( \frac{\tilde{\rho}}{\rho} \right)^4$$

$$\tilde{\rho} = \left( \frac{8}{3\pi^6} \right)^{1/4} S M_N$$

The longevity of the universe implies

$$\tilde{\rho} > 0.140 M_{\odot} R_E^{-3} = (1.5 \text{ MeV})^4$$

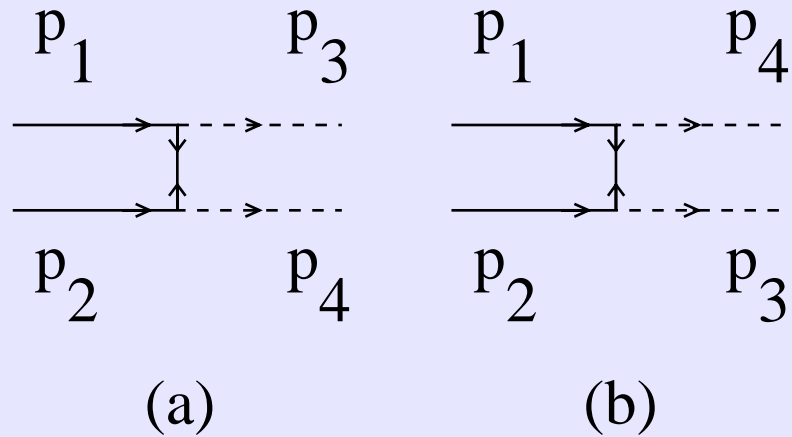
not far from the average density of the prototype white dwarf:

$$\rho_{WD} = \frac{3}{4\pi} M_{\odot} R_E^{-3}$$

transition rate:  $\frac{1}{N} \frac{dN}{dt} = AV e^{-(\tilde{\rho}/\rho)^4}$

	M	V	$\rho$	$V e^{-(\tilde{\rho}/\rho)^4}$
Cluster core	$10^{14}$	$4 \cdot 10^{44}$	$2.5 \cdot 10^{-31}$	$e^{-10^{119}}$
Sun	1	$4 \cdot 10^6$	$2.5 \cdot 10^{-7}$	$e^{-10^{23}}$
Earth	$10^{-6}$	4	$2.5 \cdot 10^{-7}$	$e^{-10^{23}}$
White Dwarf	1	4	.25	3.7
Neutron Star	1	$0.8 \cdot 10^{-9}$	$1.2 \cdot 10^9$	$0.8 \cdot 10^{-9}$
U <sup>238</sup> Nucleus	$2 \cdot 10^{-55}$	$0.8 \cdot 10^{-63}$	$2.5 \cdot 10^8$	$0.8 \cdot 10^{-63}$

Masses, volumes, mean densities, and relative transition rates for  $\tilde{\rho} = 0.14$



$$e^- e^- \rightarrow \tilde{e}^- \tilde{e}^-$$

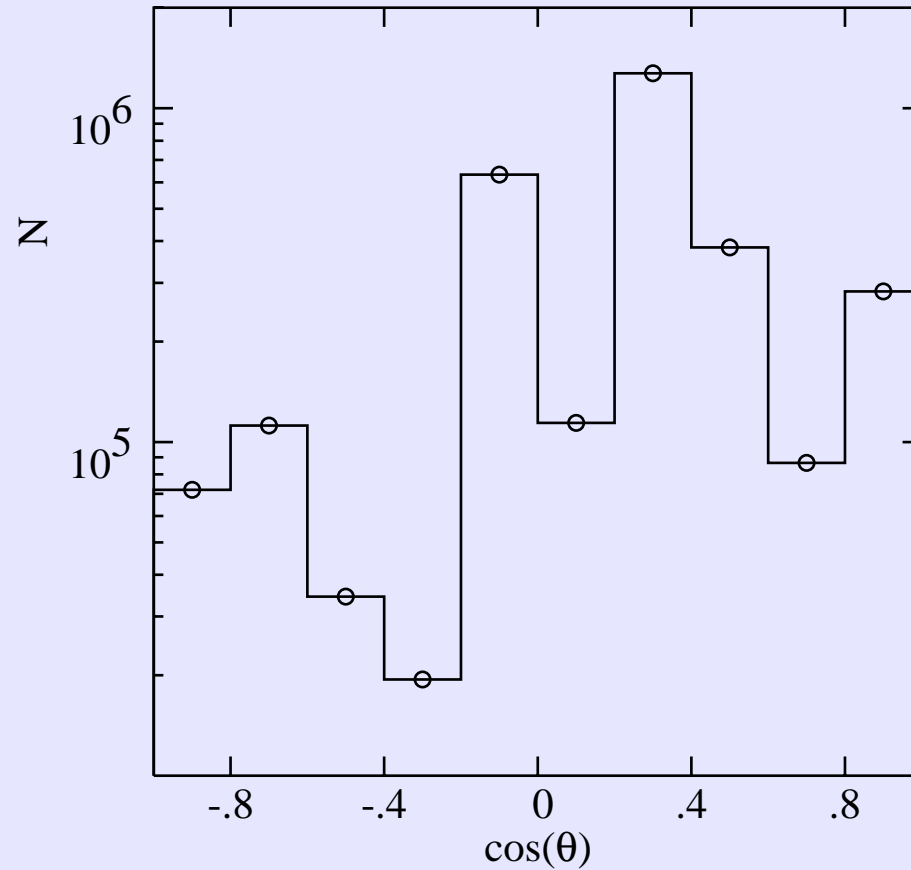
electron to selectron pair conversion via photino exchange

[LC and I. Perevalova, Phys Rev D, March 05](#)

also for quarks:  $qq \rightarrow \tilde{q}\tilde{q}$  and nucleons:  $NN \rightarrow \tilde{N}\tilde{N}$

Four types of susy nucleons :

$$N_0 = qq\bar{q} \quad , \quad N_1 = qq\tilde{q} \quad , \quad N_2 = q\tilde{q}\tilde{q} \quad , \quad N_3 = \tilde{q}\tilde{q}\tilde{q}$$



$$e^-e^- \rightarrow \tilde{e}^-\tilde{e}^-$$

angular distribution of final state selectrons  
 after 2 million events including effect of Bose enhancement

*local Fermi momentum* : 
$$p_F(r) = \left( \frac{3\pi^2 \rho(r)}{2m_N} \right)^{1/3}$$

The momentum distribution in the electron sea is no longer simply quadratic

$$\frac{dN}{dp} = \frac{4p^2}{\pi} \int_0^R r^2 dr \theta(p_F(r) - p)$$

$$p_{max} = p_F(0) = \left( \frac{3\pi^2 \rho(0)}{2M_N} \right)^{1/3}$$

<b>n</b>	<b>M</b> <b>(<math>M_{\odot}</math>)</b>	<b>R</b> <b>(<math>R_E</math>)</b>	<b><math>\rho_0</math></b> <b>(<math>\frac{M_{\odot}}{R_E^3}</math>)</b>	<b><math>\bar{\rho}</math></b> <b>(<math>\frac{M_{\odot}}{R_E^3}</math>)</b>	<b><math>\bar{E}</math></b> <b>(MeV)</b>	<b><math>E_{tot}</math></b> <b>(ergs)</b>
<b>1</b>	<b>0.271</b>	<b>1.9518</b>	<b>0.056</b>	<b>0.009</b>	<b>0.013</b>	<b><math>3.5 \cdot 10^{48}</math></b>
<b>2</b>	<b>0.512</b>	<b>1.4499</b>	<b>0.297</b>	<b>0.040</b>	<b>0.034</b>	<b><math>1.7 \cdot 10^{49}</math></b>
<b>3</b>	<b>0.811</b>	<b>1.0505</b>	<b>1.572</b>	<b>0.167</b>	<b>0.078</b>	<b><math>6.1 \cdot 10^{49}</math></b>
<b>4</b>	<b>1.072</b>	<b>0.7376</b>	<b>8.321</b>	<b>0.637</b>	<b>0.158</b>	<b><math>1.7 \cdot 10^{50}</math></b>
<b>5</b>	<b>1.242</b>	<b>0.5002</b>	<b>44.053</b>	<b>2.366</b>	<b>0.291</b>	<b><math>3.5 \cdot 10^{50}</math></b>
<b>6</b>	<b>1.330</b>	<b>0.3269</b>	<b>233.238</b>	<b>9.077</b>	<b>0.510</b>	<b><math>6.6 \cdot 10^{50}</math></b>
<b>7</b>	<b>1.369</b>	<b>0.2062</b>	<b>1234.875</b>	<b>37.212</b>	<b>0.877</b>	<b><math>1.2 \cdot 10^{51}</math></b>

**physical effects influencing the burst duration:**

- 1. The bubble growth time.**
- 2. The light crossing time.**
- 3. The free collapse time.**
- 4. Radiation pressure.**

bubble growth time assumed to be governed approximately by the speed of sound in dense matter.

$$v_s(r) = \sqrt{\frac{3P}{\rho}}$$

In a gravitationally bound medium of constant density  $v_s(r) = \sqrt{\frac{2}{3}\gamma\pi G_N\rho(R^2 - r^2)}$

$\gamma$  is the ratio of specific heats (5/3 for a monatomic gas), and  $R$  is the stellar surface radius at which the pressure vanishes.

bubble growth time for constant density:  $\tau = \int_0^R dr/v_s(r) = \frac{\pi R}{2v_s(0)}$

$$\frac{1}{N} \frac{dN}{dt} = 4\pi A \int_0^R r^2 dr e^{-\left(\frac{\tilde{p}}{\rho(r)}\right)^4}$$



**classical free collapse time after relief of Pauli blocking:**

$$\tau_c = \frac{\pi}{2} \left( \frac{8\pi G_N \rho}{3} \right)^{-1/2} \approx 1.5s \left( \frac{\rho}{\rho_{WD}} \right)^{-1/2}$$

**near Schwarzschild radius, GR effects dilate collapse time as seen by a distant observer:**

Schwarzschild radius :  $r^* = \frac{2G_N M}{c^2}$

$$r - r^* = (r_0 - r^*) e^{-(t-t_0)c/r^*}$$

$$E = E_{em} \frac{e^{-(t-t_0)c/2r^*}}{\sqrt{(1 + e^{-(t-t_0)c/r^*})}}$$

$$\frac{dN}{dt} = N(M)4\pi A \int_0^R r^2 dr e^{-(\tilde{\rho}/\rho(r))^4}$$

*bubble growth time :*  $\tau_0 = \int_0^R dr/v_s(r)$

$\bar{\tau}$  = bubble growth time assuming constant (average) density

$\tau_c$  = classical free collapse time

n	M ( $M_\odot$ )	R ( $R_E$ )	$\tau_0$ (s)	$\bar{\tau}$ (s)	$\tau_c$ (s)	$E_{tot}$ (ergs)	$\tilde{\rho} = 0.14$	$\frac{1}{A} \frac{dN}{dt}$ $\tilde{\rho} = 0.7$	$\tilde{\rho} = 3.5$
1	.271	1.92	1.42	12.3	7.82	$3.5 \cdot 10^{48}$	$10^{-11}$	0	0
2	.512	1.44	0.69	5.80	3.70	$1.7 \cdot 10^{49}$	<u>18</u>	$10^{-7}$	$4 \cdot 10^{-44}$
3	.811	1.05	0.35	2.86	1.82	$6.1 \cdot 10^{49}$	0.78	12	$7.9 \cdot 10^{-8}$
4	1.07	.737	0.18	1.47	0.93	$1.7 \cdot 10^{50}$	0.46	<u>20</u>	31
5	1.24	.500	.095	0.76	0.49	$3.5 \cdot 10^{50}$	0.19	11	<u>44</u>
6	1.33	.327	.049	0.39	0.25	$6.6 \cdot 10^{50}$	0.062	4.4	24
7	1.37	.206	0.024	0.19	0.12	$1.2 \cdot 10^{51}$	0.017	1.4	9.2

## summary

The susy star model of gamma ray bursts provides a new framework in which to discuss violent astrophysical events

### conceptual advantages:

- 1.) Connection to vacuum decay and string landscape
- 2.) Adequate explosive energy release
- 3.) jet seeding mechanism through Bose enhancement
- 4.) possible solution to the baryon loading problem

### zeroth order quantitative predictions:

- 1.) roughly correct mean photon energy
- 2.) roughly correct total energy release
- 3.) roughly correct burst duration

taking density inhomogeneity into account, the average photon energy from a susy phase transition is several times greater than that assuming constant density.

### still to be understood:

- 1.) spikey behavior of bursts
- 2.) relation to supernovae
- 3.) bi-modal duration distribution
- 4.) identifiable bursts seem to be clumped at red-shift  $z \approx 1$